

Well-posedness of the linearized plasma-vacuum interface problem

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We consider the free boundary problem for the plasma-vacuum interface in ideal compressible magnetohydrodynamics (MHD). In the plasma region the flow is governed by the usual compressible MHD equations, while in the vacuum region we consider the *pre-Maxwell dynamics* for the magnetic field. At the free-interface we assume that the total pressure is continuous and that the magnetic field is tangent to the boundary. The plasma density does not go to zero continuously at the interface, but has a jump, meaning that it is bounded away from zero in the plasma region and it is identically zero in the vacuum region. Under a suitable stability condition satisfied at each point of the plasma-vacuum interface, we prove the well-posedness of the linearized problem in conormal Sobolev spaces.

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1. Introduction

Consider the equations of ideal compressible MHD:

$$\begin{cases} \partial_t \rho + \operatorname{div}(\rho v) = 0, \\ \partial_t(\rho v) + \operatorname{div}(\rho v \otimes v - H \otimes H) + \nabla q = 0, \\ \partial_t H - \nabla \times (v \times H) = 0, \\ \partial_t(\rho e + \frac{1}{2}|H|^2) + \operatorname{div}((\rho e + p)v + H \times (v \times H)) = 0, \end{cases} \quad (1.1)$$

where ρ denotes density, $v \in \mathbb{R}^3$ plasma velocity, $H \in \mathbb{R}^3$ magnetic field, $p = p(\rho, S)$ pressure, $q = p + \frac{1}{2}|H|^2$ total pressure, S entropy, $e = E + \frac{1}{2}|v|^2$ total energy, and $E = E(\rho, S)$ internal energy. With a state equation of gas, $\rho = \rho(p, S)$, and the first principle of thermodynamics, (1.1) is a closed system.

System (1.1) is supplemented by the divergence constraint

$$\operatorname{div} H = 0 \quad (1.2)$$

on the initial data. As is known, taking into account (1.2), we can easily symmetrize system (1.1) by rewriting it in the nonconservative form

$$\begin{cases} \frac{\rho_p}{\rho} \frac{dp}{dt} + \operatorname{div} v = 0, & \rho \frac{dv}{dt} - (H \cdot \nabla)H + \nabla q = 0, \\ \frac{dH}{dt} - (H \cdot \nabla)v + H \operatorname{div} v = 0, & \frac{dS}{dt} = 0, \end{cases} \quad (1.3)$$

where $\rho_p \equiv \partial \rho / \partial p$ and $d/dt = \partial_t + (v \cdot \nabla)$. A different symmetrization is obtained if we consider q instead of p . In terms of q the equation for the pressure in (1.3) takes the form

$$\frac{\rho_p}{\rho} \left\{ \frac{dq}{dt} - H \cdot \frac{dH}{dt} \right\} + \operatorname{div} v = 0, \quad (1.4)$$

where it is understood that now $\rho = \rho(q - |H|^2/2, S)$ and similarly for ρ_p . Then we derive $\operatorname{div} v$ from (1.4) and rewrite the equation for the magnetic field in (1.3) as

$$\frac{dH}{dt} - (H \cdot \nabla)v - \frac{\rho_p}{\rho} H \left\{ \frac{dq}{dt} - H \cdot \frac{dH}{dt} \right\} = 0. \quad (1.5)$$

Substituting (1.4), (1.5) in (1.3) then gives the following symmetric system

$$\begin{aligned} & \begin{pmatrix} \rho_p/\rho & \underline{0} & -(\rho_p/\rho)H & 0 \\ \underline{0}^T & \rho I_3 & 0_3 & \underline{0}^T \\ -(\rho_p/\rho)H^T & 0_3 & I_3 + (\rho_p/\rho)H \otimes H & \underline{0}^T \\ 0 & \underline{0} & \underline{0} & 1 \end{pmatrix} \partial_t \begin{pmatrix} q \\ v \\ H \\ S \end{pmatrix} + \\ & + \begin{pmatrix} (\rho_p/\rho)v \cdot \nabla & \nabla \cdot & -(\rho_p/\rho)Hv \cdot \nabla & 0 \\ \nabla & \rho v \cdot \nabla I_3 & -H \cdot \nabla I_3 & \underline{0}^T \\ -(\rho_p/\rho)H^T v \cdot \nabla & -H \cdot \nabla I_3 & (I_3 + (\rho_p/\rho)H \otimes H)v \cdot \nabla & \underline{0}^T \\ 0 & \underline{0} & \underline{0} & v \cdot \nabla \end{pmatrix} \begin{pmatrix} q \\ v \\ H \\ S \end{pmatrix} = 0, \end{aligned} \quad (1.6)$$

where $\underline{0} = (0, 0, 0)$. Given this symmetrization, as the unknown we can choose the vector $U = U(t, x) = (q, v, H, S)$. For the sake of brevity we write system (1.6) in the form

$$A_0(U) \partial_t U + \sum_{j=1}^3 A_j(U) \partial_j U = 0, \quad (1.7)$$

which is symmetric hyperbolic provided the hyperbolicity condition $A_0 > 0$ holds:

$$\rho > 0, \quad \rho_p > 0. \quad (1.8)$$

Plasma-vacuum interface problems for system (1.1) appear in the mathematical modeling of plasma confinement by magnetic fields (see, e.g., [10]). In this model the plasma is confined inside a perfectly conducting rigid wall and isolated from it by a vacuum region, due to the effect of strong magnetic fields. This subject is very popular since the 1950–70's, but most of theoretical studies are devoted to finding stability criteria of equilibrium states. The typical work in this direction is

the classical paper of Bernstein et al. [3]. In astrophysics, the plasma-vacuum interface problem can be used for modeling the motion of a star or the solar corona when magnetic fields are taken into account.

According to our knowledge there are still no well-posedness results for full (*non-stationary*) plasma-vacuum models. More precisely, an energy a priori estimate in Sobolev spaces for the linearization of a plasma-vacuum interface problem (see its description just below) was proved in [23], but the existence of solutions to this problem remained open. In fact, the proof of existence of solutions is the main goal of the present paper.

Let $\Omega^+(t)$ and $\Omega^-(t)$ be space-time domains occupied by the plasma and the vacuum respectively. That is, in the domain $\Omega^+(t)$ we consider system (1.1) (or (1.7)) governing the motion of an ideal plasma and in the domain $\Omega^-(t)$ we have the elliptic (div-curl) system

$$\nabla \times \mathcal{H} = 0, \quad \operatorname{div} \mathcal{H} = 0, \quad (1.9)$$

describing the vacuum magnetic field $\mathcal{H} \in \mathbb{R}^3$. Here, as in [3, 10], we consider so-called *pre-Maxwell dynamics*. That is, as usual in nonrelativistic MHD, we neglect the displacement current $(1/c) \partial_t E$, where c is the speed of light and E is the electric field.

Let us assume that the interface between plasma and vacuum is given by a hypersurface $\Gamma(t) = \{F(t, x) = 0\}$. It is to be determined and moves with the velocity of plasma particles at the boundary:

$$\frac{dF}{dt} = 0 \quad \text{on } \Gamma(t) \quad (1.10)$$

(for all $t \in [0, T]$). As F is an unknown of the problem, this is a free-boundary problem. The plasma variable U is connected with the vacuum magnetic field \mathcal{H} through the relations [3, 10]

$$[q] = 0, \quad H \cdot N = 0, \quad \mathcal{H} \cdot N = 0, \quad \text{on } \Gamma(t), \quad (1.11)$$

where $N = \nabla F$ and $[q] = q|_F - \frac{1}{2}|\mathcal{H}|_F^2$ denotes the jump of the total pressure across the interface. These relations together with (1.10) are the boundary conditions at the interface $\Gamma(t)$.

As in [12, 22], we will assume that for problem (1.1), (1.9)–(1.11) the hyperbolicity conditions (1.8) are assumed to be satisfied in $\Omega^+(t)$ up to the boundary $\Gamma(t)$, i.e., the plasma density does not go to zero continuously, but has a jump (clearly in the vacuum region $\Omega^-(t)$ the density is identically zero). This assumption is compatible with the continuity of the total pressure in (1.11). For instance, in the case of ideal polytropic gases one has $p = A\rho^\gamma e^S$ with $A > 0, \gamma > 1$. Then the continuity of the total pressure at Γ requires $(A\rho^\gamma e^S + \frac{1}{2}|H|^2)|_{\Gamma_+} - \frac{1}{2}|\mathcal{H}|^2 = 0|_{\Gamma_-}$, which may be obtained also for densities ρ discontinuous across Γ . Differently, in the absence of the magnetic field, the continuity of the pressure yields the continuity of the density so that the boundary condition becomes $\rho|_{\Gamma_+} = 0$.

Since the interface moves with the velocity of plasma particles at the boundary, by introducing the Lagrangian coordinates one can reduce the original problem to that in a fixed domain. This approach has been recently employed with success in a series of papers on the Euler equations in vacuum, see [6–9, 12]. However, as, for example, for contact discontinuities in various models of fluid dynamics (e.g., for current-vortex sheets [4, 21]), this approach seems hardly applicable for problem (1.1), (1.9)–(1.11). Therefore, we will work in the Eulerian coordinates and for technical simplicity we will assume that the space-time domains $\Omega^\pm(t)$ have the following form.

Let us assume that the moving interface $\Gamma(t)$ takes the form

$$\Gamma(t) := \{(x_1, x') \in \mathbb{R}^3, x_1 = \varphi(t, x')\},$$

where $t \in [0, T]$ and $x' = (x_2, x_3)$. Then we have $\Omega^\pm(t) = \{x_1 \gtrless \varphi(t, x')\}$. With our parametrization of $\Gamma(t)$, an equivalent formulation of the boundary conditions (1.10), (1.11) at the interface is

$$\partial_t \varphi = v_N, \quad [q] = 0, \quad H_N = 0, \quad \mathcal{H}_N = 0 \quad \text{on } \Gamma(t), \quad (1.12)$$

where $v_N = v \cdot N$, $H_N = H \cdot N$, $\mathcal{H}_N = \mathcal{H} \cdot N$, $N = (1, -\partial_2 \varphi, -\partial_3 \varphi)$.

System (1.7), (1.9), (1.12) is supplemented with initial conditions

$$\begin{aligned} U(0, x) &= U_0(x), & x &\in \Omega^+(0), & \varphi(0, x) &= \varphi_0(x), & x &\in \Gamma, \\ \mathcal{H}(0, x) &= \mathcal{H}_0(x), & x &\in \Omega^-(0), \end{aligned} \quad (1.13)$$

From the mathematical point of view, a natural wish is to find conditions on the initial data providing the existence and uniqueness on some time interval $[0, T]$ of a solution $(U, \mathcal{H}, \varphi)$ to problem (1.7), (1.9), (1.12), (1.13) in Sobolev spaces. Since (1.1) is a system of hyperbolic conservation laws that can produce shock waves and other types of strong discontinuities (e.g., current-vortex sheets [21]), it is natural to expect to obtain only local-in-time existence theorems. Notice that (1.7), (1.9) is a coupled hyperbolic-elliptic system.

We must regard the boundary conditions on H in (1.12) as the restriction on the initial data (1.13). More precisely, we can prove that a solution of (1.7), (1.12) (if it exists for all $t \in [0, T]$) satisfies

$$\operatorname{div} H = 0 \quad \text{in } \Omega^\pm(t) \quad \text{and} \quad H_N = 0 \quad \text{on } \Gamma(t),$$

for all $t \in [0, T]$, if the latter were satisfied at $t = 0$, i.e., for the initial data (1.13). In particular, the fulfillment of $\operatorname{div} H = 0$ implies that systems (1.1) and (1.7) are equivalent on solutions of problem (1.7), (1.12), (1.13).

The remainder of the paper is organized as follows. In the next section we introduce an equivalent formulation in the fixed domain with flat boundary. In Section 2 we formulate the linearized problem associated to (1.17)–(1.19) and introduce suitable decompositions of the magnetic fields to reduce it to that with homogeneous boundary conditions and homogeneous linearized “vacuum” equations. In fact, for proving the basic a priori energy estimate, it is convenient to have the vacuum magnetic field satisfying homogeneous equations and boundary conditions as in (2.14), and the plasma magnetic field satisfying homogeneous constraints (2.23), (2.24). Thus we introduce the decomposition $\mathcal{H} = \mathcal{H}' + \mathcal{H}''$ in the vacuum side, with \mathcal{H}' solution of (2.14), and \mathcal{H}'' taking all the nonhomogeneous part (2.13), and the decomposition (2.21) in the plasma side.

The main result of the paper is stated in Section 4. We prove the existence of a unique solution to the linearized hyperbolic-elliptic problem (2.29) satisfying the a priori estimate (4.2). The a priori estimate (4.2) improves the similar estimate firstly proved in [23].

In Section 5 we introduce a fully hyperbolic regularization (5.1) of the coupled hyperbolic-elliptic system (2.29). In Section 6 we show an a priori estimate of solutions uniform in the small parameter ε of regularization. In Section 7 we prove the existence of the solution to the hyperbolic regularizing problem (5.1). For it one main difficulty is the fact that the problem is non standard, due to the coupling with the front, and that the Kreiss-Lopatinskii condition doesn’t hold uniformly, so that for instance the approach of [2] does not apply. Other difficulties are due to the characteristic boundary and the lack of *reflexivity* (in the sense of Ohkubo [17], [2]). In Section 7 we find a relatively simple proof of existence by means of an alternative formulation and a fixed point argument. In Section 8 we conclude the proof of Theorem 4.1 by passing to the limit as $\varepsilon \rightarrow 0$. Sections 9, 10, 11 are devoted to the proof of some technical results.

1.1 An equivalent formulation in the fixed domain

Let us denote

$$\Omega^\pm := \mathbb{R}^3 \cap \{x_1 \gtrless 0\}, \quad \Gamma := \mathbb{R}^3 \cap \{x_1 = 0\}.$$

We want to reduce the free boundary problem (1.7), (1.9), (1.12), (1.13) to the fixed domains Ω^\pm . For this purpose we introduce a suitable change of variables that is inspired by Lannes [11]. In all what follows, $H^s(\omega)$ denotes the Sobolev space of order s on a domain ω . The following lemma shows how to lift functions from Γ to \mathbb{R}^3 . An important point is the regularization of one half derivative of the lifting function Ψ w.r.t. the given function φ . For instance, there is no such regularization in the lifting function chosen in [13, 14].

LEMMA 1.1 Let $m \geq 3$. For all $\epsilon > 0$ there exists a continuous linear map $\varphi \in H^{m-0.5}(\mathbb{R}^2) \mapsto \Psi \in H^m(\mathbb{R}^3)$ such that $\Psi(0, x') = \varphi(x')$, $\partial_1 \Psi(0, x') = 0$ on Γ , and

$$\|\partial_1 \Psi\|_{L^\infty(\mathbb{R}^3)} \leq \epsilon \|\varphi\|_{H^2(\mathbb{R}^2)}. \quad (1.14)$$

We give the proof of Lemma 1.1 in Section 10 at the end of this article. The following lemma gives the time-dependent version of Lemma 1.1.

LEMMA 1.2 Let $m \geq 3$ be an integer and let $T > 0$. For all $\epsilon > 0$ there exists a continuous linear map $\varphi \in \cap_{j=0}^{m-1} \mathcal{C}^j([0, T]; H^{m-j-0.5}(\mathbb{R}^2)) \mapsto \Psi \in \cap_{j=0}^{m-1} \mathcal{C}^j([0, T]; H^{m-j}(\mathbb{R}^3))$ such that $\Psi(t, 0, x') = \varphi(t, x')$, $\partial_1 \Psi(t, 0, x') = 0$ on Γ , and

$$\|\partial_1 \Psi\|_{\mathcal{C}([0, T]; L^\infty(\mathbb{R}^3))} \leq \epsilon \|\varphi\|_{\mathcal{C}([0, T]; H^2(\mathbb{R}^2))}. \quad (1.15)$$

Furthermore, there exists a constant $C > 0$ that is independent of T and only depends on m , such that

$$\forall \varphi \in \cap_{j=0}^{m-1} \mathcal{C}^j([0, T]; H^{m-j-0.5}(\mathbb{R}^2)), \quad \forall j = 0, \dots, m-1, \quad \forall t \in [0, T], \\ \|\partial_t^j \Psi(t, \cdot)\|_{H^{m-j}(\mathbb{R}^3)} \leq C \|\partial_t^j \varphi(t, \cdot)\|_{H^{m-j-0.5}(\mathbb{R}^2)}.$$

The proof of Lemma 1.2 is also postponed to Section 10. The diffeomorphism that reduces the free boundary problem (1.7), (1.12), (1.13) to the fixed domains Ω^\pm is given in the following lemma.

LEMMA 1.3 Let $m \geq 3$ be an integer. For all $T > 0$, and for all $\varphi \in \cap_{j=0}^{m-1} \mathcal{C}^j([0, T]; H^{m-j-0.5}(\mathbb{R}^2))$, satisfying without loss of generality $\|\varphi\|_{\mathcal{C}([0, T]; H^2(\mathbb{R}^2))} \leq 1$, there exists a function $\Psi \in \cap_{j=0}^{m-1} \mathcal{C}^j([0, T]; H^{m-j}(\mathbb{R}^3))$ such that the function

$$\Phi(t, x) := (x_1 + \Psi(t, x), x'), \quad (t, x) \in [0, T] \times \mathbb{R}^3, \quad (1.16)$$

defines an H^m -diffeomorphism of \mathbb{R}^3 for all $t \in [0, T]$. Moreover, there holds $\partial_t^j (\Phi - Id) \in \mathcal{C}([0, T]; H^{m-j}(\mathbb{R}^3))$ for $j = 0, \dots, m-1$, $\Phi(t, 0, x') = (\varphi(t, x'), x')$, $\partial_1 \Phi(t, 0, x') = (1, 0, 0)$.

Proof of Lemma 1.3. The proof follows directly from Lemma 1.2 because

$$\partial_1 \Phi(t, x) = 1 + \partial_1 \Psi(t, x) \geq 1 - \|\partial_1 \Psi(t, \cdot)\|_{\mathcal{C}([0, T]; L^\infty(\mathbb{R}^3))} \geq 1 - \epsilon \|\varphi\|_{\mathcal{C}([0, T]; H^2(\mathbb{R}^2))} \geq 1/2,$$

provided ϵ is taken sufficiently small, e.g. $\epsilon < 1/2$. The other properties of Φ follow directly from Lemma 1.2. \square

158 We introduce the change of independent variables defined by (1.16) by setting

$$159 \quad \tilde{U}(t, x) := U(t, \Phi(t, x)), \quad \tilde{\mathcal{H}}(t, x) := \mathcal{H}(t, \Phi(t, x)).$$

Dropping for convenience tildes in \tilde{U} and $\tilde{\mathcal{H}}$, problem (1.7), (1.9)–(1.12), (1.13) can be reformulated on the fixed reference domains Ω^\pm as

$$\mathbb{P}(U, \Psi) = 0 \quad \text{in } [0, T] \times \Omega^+, \quad \mathbb{V}(\mathcal{H}, \Psi) = 0 \quad \text{in } [0, T] \times \Omega^-, \quad (1.17)$$

$$\mathbb{B}(U, \mathcal{H}, \varphi) = 0 \quad \text{on } [0, T] \times \Gamma, \quad (1.18)$$

$$(U, \mathcal{H})|_{t=0} = (U_0, \mathcal{H}_0) \quad \text{in } \Omega^+ \times \Omega^-, \quad \varphi|_{t=0} = \varphi_0 \quad \text{on } \Gamma, \quad (1.19)$$

where $\mathbb{P}(U, \Psi) = P(U, \Psi)U$,

$$\begin{aligned} P(U, \Psi) &= A_0(U)\partial_t + \tilde{A}_1(U, \Psi)\partial_1 + A_2(U)\partial_2 + A_3(U)\partial_3, \\ \tilde{A}_1(U, \Psi) &= \frac{1}{\partial_1 \Phi_1} \left(A_1(U) - A_0(U)\partial_t \Psi - \sum_{k=2}^3 A_k(U)\partial_k \Psi \right), \\ \mathbb{V}(\mathcal{H}, \Psi) &= \begin{pmatrix} \nabla \times \mathfrak{H} \\ \operatorname{div} \mathfrak{h} \end{pmatrix}, \\ \mathfrak{H} &= (\mathcal{H}_1 \partial_1 \Phi_1, \mathcal{H}_{\tau_2}, \mathcal{H}_{\tau_3}), \quad \mathfrak{h} = (\mathcal{H}_N, \mathcal{H}_2 \partial_1 \Phi_1, \mathcal{H}_3 \partial_1 \Phi_1), \\ \mathcal{H}_N &= \mathcal{H}_1 - \mathcal{H}_2 \partial_2 \Psi - \mathcal{H}_3 \partial_3 \Psi, \quad \mathcal{H}_{\tau_i} = \mathcal{H}_1 \partial_i \Psi + \mathcal{H}_i, \quad i = 2, 3, \\ \mathbb{B}(U, \mathcal{H}, \varphi) &= \begin{pmatrix} \partial_t \varphi - v_N|_{x_1=0} \\ [q] \\ \mathcal{H}_{N|x_1=0} \end{pmatrix}, \quad [q] = q|_{x_1=0} - \frac{1}{2} |\mathcal{H}|_{x_1=0}^2, \\ v_N &= v_1 - v_2 \partial_2 \Psi - v_3 \partial_3 \Psi. \end{aligned}$$

160 To avoid an overload of notation we have denoted by the same symbols v_N, \mathcal{H}_N here above and
161 v_N, \mathcal{H}_N as in (1.12). Notice that $v_N|_{x_1=0} = v_1 - v_2 \partial_2 \varphi - v_3 \partial_3 \varphi$, $\mathcal{H}_N|_{x_1=0} = \mathcal{H}_1 - \mathcal{H}_2 \partial_2 \varphi - \mathcal{H}_3 \partial_3 \varphi$,
162 as in the previous definition in (1.12).

163 We did not include in problem (1.17)–(1.19) the equation

$$164 \quad \operatorname{div} h = 0 \quad \text{in } [0, T] \times \Omega^+, \quad (1.20)$$

165 and the boundary condition

$$166 \quad H_N = 0 \quad \text{on } [0, T] \times \Gamma, \quad (1.21)$$

167 where $h = (H_N, H_2 \partial_1 \Phi_1, H_3 \partial_1 \Phi_1)$, $H_N = H_1 - H_2 \partial_2 \Psi - H_3 \partial_3 \Psi$, because they are just
168 restrictions on the initial data (1.19). More precisely, referring to [21] for the proof, we have the
169 following proposition.

170 **PROPOSITION 1.4** Let the initial data (1.19) satisfy (1.20) and (1.21) for $t = 0$. If $(U, \mathcal{H}, \varphi)$ is a
171 solution of problem (1.17)–(1.19), then this solution satisfies (1.20) and (1.21) for all $t \in [0, T]$.

172 Note that Proposition 1.4 stays valid if in (1.17) we replace system $\mathbb{P}(U, \Psi) = 0$ by system
173 (1.1) in the straightened variables. This means that these systems are equivalent on solutions of our
174 plasma-vacuum interface problem and we may justifiably replace the conservation laws (1.1) by
175 their nonconservative form (1.7).

2. The linearized problem

2.1 Basic state

Let us denote

$$Q_T^\pm := (-\infty, T] \times \Omega^\pm, \quad \omega_T := (-\infty, T] \times \Gamma.$$

Let

$$(\widehat{U}(t, x), \widehat{\mathcal{H}}(t, x), \widehat{\varphi}(t, x')) \quad (2.1)$$

be a given sufficiently smooth vector-function with $\widehat{U} = (\widehat{q}, \widehat{v}, \widehat{H}, \widehat{S})$, respectively defined on Q_T^+, Q_T^-, ω_T , with

$$\begin{aligned} \|\widehat{U}\|_{W^{2,\infty}(Q_T^+)} + \|\partial_1 \widehat{U}\|_{W^{2,\infty}(Q_T^+)} + \|\widehat{\mathcal{H}}\|_{W^{2,\infty}(Q_T^-)} + \|\widehat{\varphi}\|_{W^{3,\infty}([0,T] \times \mathbb{R}^2)} &\leq K, \\ \|\widehat{\varphi}\|_{C([0,T]; H^2(\mathbb{R}^2))} &\leq 1, \end{aligned} \quad (2.2)$$

where $K > 0$ is a constant. Corresponding to the given $\widehat{\varphi}$ we construct $\widehat{\psi}$ and the diffeomorphism $\widehat{\Phi}$ as in Lemmata 1.2 and 1.3 such that

$$\partial_1 \widehat{\Phi}_1 \geq 1/2.$$

We assume that the basic state (2.1) satisfies (for some positive $\rho_0, \rho_1 \in \mathbb{R}$)

$$\rho(\widehat{p}, \widehat{S}) \geq \rho_0 > 0, \quad \rho_p(\widehat{p}, \widehat{S}) \geq \rho_1 > 0 \quad \text{in } \overline{Q_T^+}, \quad (2.3)$$

$$\partial_t \widehat{H} + \frac{1}{\partial_1 \widehat{\Phi}_1} \left\{ (\widehat{w} \cdot \nabla) \widehat{H} - (\widehat{h} \cdot \nabla) \widehat{v} + \widehat{H} \operatorname{div} \widehat{u} \right\} = 0 \quad \text{in } Q_T^+, \quad (2.4)$$

$$\operatorname{div} \widehat{h} = 0 \quad \text{in } Q_T^-, \quad (2.5)$$

$$\partial_t \widehat{\varphi} - \widehat{v}_N = 0, \quad \widehat{\mathcal{H}}_N = 0 \quad \text{on } \omega_T, \quad (2.6)$$

where all the “hat” values are determined like corresponding values for $(U, \mathcal{H}, \varphi)$, i.e.,

$$\begin{aligned} \widehat{\mathfrak{h}} &= (\widehat{\mathcal{H}}_1 \partial_1 \widehat{\Phi}_1, \widehat{\mathcal{H}}_{\tau_2}, \widehat{\mathcal{H}}_{\tau_3}), \quad \widehat{h} = (\widehat{\mathcal{H}}_N, \widehat{\mathcal{H}}_2 \partial_1 \widehat{\Phi}_1, \widehat{\mathcal{H}}_3 \partial_1 \widehat{\Phi}_1), \quad \widehat{h} = (\widehat{H}_N, \widehat{H}_2 \partial_1 \widehat{\Phi}_1, \widehat{H}_3 \partial_1 \widehat{\Phi}_1), \\ \widehat{p} &= \widehat{q} - |\widehat{H}|^2/2, \quad \widehat{v}_N = \widehat{v}_1 - \widehat{v}_2 \partial_2 \widehat{\psi} - \widehat{v}_3 \partial_3 \widehat{\psi}, \quad \widehat{\mathcal{H}}_N = \widehat{\mathcal{H}}_1 - \widehat{\mathcal{H}}_2 \partial_2 \widehat{\psi} - \widehat{\mathcal{H}}_3 \partial_3 \widehat{\psi}, \end{aligned}$$

and where

$$\widehat{u} = (\widehat{v}_N, \widehat{v}_2 \partial_1 \widehat{\Phi}_1, \widehat{v}_3 \partial_1 \widehat{\Phi}_1), \quad \widehat{w} = \widehat{u} - (\partial_t \widehat{\psi}, 0, 0).$$

Note that (2.2) yields

$$\|\nabla_{t,x} \widehat{\psi}\|_{W^{2,\infty}([0,T] \times \mathbb{R}^3)} \leq C(K),$$

where $\nabla_{t,x} = (\partial_t, \nabla)$ and $C = C(K) > 0$ is a constant depending on K .

It follows from (2.4) that the constraints

$$\operatorname{div} \widehat{h} = 0 \quad \text{in } Q_T^+, \quad \widehat{H}_N = 0 \quad \text{on } \omega_T, \quad (2.7)$$

are satisfied for the basic state (2.1) if they hold at $t = 0$ (see [21] for the proof). Thus, for the basic state we also require the fulfillment of conditions (2.7) at $t = 0$.

2.2 Linearized problem

The linearized equations for (1.17), (1.18) read:

$$\begin{aligned}\mathbb{P}'(\widehat{U}, \widehat{\Psi})(\delta U, \delta \Psi) &:= \frac{d}{d\varepsilon} \mathbb{P}(U_\varepsilon, \Psi_\varepsilon)|_{\varepsilon=0} = f \quad \text{in } Q_T^+, \\ \mathbb{V}'(\widehat{\mathcal{H}}, \widehat{\Psi})(\delta \mathcal{H}, \delta \Psi) &:= \frac{d}{d\varepsilon} \mathbb{V}(\mathcal{H}_\varepsilon, \Psi_\varepsilon)|_{\varepsilon=0} = \mathcal{G}' \quad \text{in } Q_T^-, \\ \mathbb{B}'(\widehat{U}, \widehat{\mathcal{H}}, \widehat{\varphi})(\delta U, \delta \mathcal{H}, \delta \varphi) &:= \frac{d}{d\varepsilon} \mathbb{B}(U_\varepsilon, \mathcal{H}_\varepsilon, \varphi_\varepsilon)|_{\varepsilon=0} = g \quad \text{on } \omega_T,\end{aligned}$$

where $U_\varepsilon = \widehat{U} + \varepsilon \delta U$, $\mathcal{H}_\varepsilon = \widehat{\mathcal{H}} + \varepsilon \delta \mathcal{H}$, $\varphi_\varepsilon = \widehat{\varphi} + \varepsilon \delta \varphi$; $\delta \Psi$ is constructed from $\delta \varphi$ as in Lemma 1.2 and $\Psi_\varepsilon = \widehat{\Psi} + \varepsilon \delta \Psi$.

Here we introduce the source terms $f = (f_1, \dots, f_8)$, $\mathcal{G}' = (\chi, \Xi)$, $\chi = (\chi_1, \chi_2, \chi_3)$, and $g = (g_1, g_2, g_3)$ to make the interior equations and the boundary conditions inhomogeneous.

We compute the exact form of the linearized equations (below we drop δ):

$$\begin{aligned}\mathbb{P}'(\widehat{U}, \widehat{\Psi})(U, \Psi) &= P(\widehat{U}, \widehat{\Psi})U + \mathcal{C}(\widehat{U}, \widehat{\Psi})U - \{P(\widehat{U}, \widehat{\Psi})\Psi\} \frac{\partial_1 \widehat{U}}{\partial_1 \widehat{\Phi}_1} = f, \\ \mathbb{V}'(\widehat{\mathcal{H}}, \widehat{\Psi})(\mathcal{H}, \Psi) &= \mathbb{V}(\mathcal{H}, \widehat{\Psi}) + \begin{pmatrix} \nabla \widehat{\mathcal{H}}_1 \times \nabla \Psi \\ \nabla \times \begin{pmatrix} 0 \\ -\widehat{\mathcal{H}}_3 \\ \widehat{\mathcal{H}}_2 \end{pmatrix} \cdot \nabla \Psi \end{pmatrix} = \mathcal{G}', \\ \mathbb{B}'(\widehat{U}, \widehat{\mathcal{H}}, \widehat{\varphi})(U, \mathcal{H}, \varphi) &= \begin{pmatrix} \partial_t \varphi + \widehat{v}_2 \partial_2 \varphi + \widehat{v}_3 \partial_3 \varphi - v_N \\ q - \widehat{\mathcal{H}} \cdot \mathcal{H} \\ \mathcal{H}_N - \widehat{\mathcal{H}}_2 \partial_2 \varphi - \widehat{\mathcal{H}}_3 \partial_3 \varphi \end{pmatrix}_{|x_1=0} = g,\end{aligned}$$

where $q := p + \widehat{H} \cdot H$, $v_N := v_1 - v_2 \partial_2 \widehat{\Psi} - v_3 \partial_3 \widehat{\Psi}$, and the matrix $\mathcal{C}(\widehat{U}, \widehat{\Psi})$ is determined as follows:

$$\begin{aligned}\mathcal{C}(\widehat{U}, \widehat{\Psi})Y &= (Y, \nabla_y A_0(\widehat{U})) \partial_t \widehat{U} + (Y, \nabla_y \widetilde{A}_1(\widehat{U}, \widehat{\Psi})) \partial_1 \widehat{U} \\ &\quad + (Y, \nabla_y A_2(\widehat{U})) \partial_2 \widehat{U} + (Y, \nabla_y A_3(\widehat{U})) \partial_3 \widehat{U}, \\ (Y, \nabla_y A(\widehat{U})) &:= \sum_{i=1}^8 y_i \left(\frac{\partial A(Y)}{\partial y_i} \Big|_{Y=\widehat{U}} \right), \quad Y = (y_1, \dots, y_8).\end{aligned}$$

Since the differential operators $\mathbb{P}'(\widehat{U}, \widehat{\Psi})$ and $\mathbb{V}'(\widehat{\mathcal{H}}, \widehat{\Psi})$ are first-order operators in Ψ , as in [1] the linearized problem is rewritten in terms of the “good unknown”

$$\dot{U} := U - \frac{\Psi}{\partial_1 \widehat{\Phi}_1} \partial_1 \widehat{U}, \quad \dot{\mathcal{H}} := \mathcal{H} - \frac{\Psi}{\partial_1 \widehat{\Phi}_1} \partial_1 \widehat{\mathcal{H}}. \quad (2.8)$$

Taking into account assumptions (2.6) and omitting detailed calculations, we rewrite our linearized

equations in terms of the new unknowns (2.8):

$$\begin{aligned} P(\widehat{U}, \widehat{\Psi})\dot{U} + \mathcal{C}(\widehat{U}, \widehat{\Psi})\dot{U} + \frac{\Psi}{\partial_1 \widehat{\Phi}_1} \partial_1 \{\mathbb{P}(\widehat{U}, \widehat{\Psi})\} &= f, \\ \mathbb{V}(\dot{\mathcal{H}}, \widehat{\Psi}) + \frac{\Psi}{\partial_1 \widehat{\Phi}_1} \partial_1 \{\mathbb{V}(\dot{\mathcal{H}}, \widehat{\Psi})\} &= g'. \end{aligned} \quad (2.9)$$

$$\begin{aligned} \mathbb{B}'(\widehat{U}, \dot{\mathcal{H}}, \widehat{\varphi})(\dot{U}, \dot{\mathcal{H}}, \varphi) &:= \mathbb{B}'(\widehat{U}, \dot{\mathcal{H}}, \widehat{\varphi})(U, \mathcal{H}, \varphi) \\ &= \begin{pmatrix} \partial_t \varphi + \widehat{v}_2 \partial_2 \varphi + \widehat{v}_3 \partial_3 \varphi - \widehat{v}_N - \varphi \partial_1 \widehat{v}_N \\ \dot{q} - \dot{\mathcal{H}} \cdot \dot{\mathcal{H}} + [\partial_1 \dot{q}] \varphi \\ \dot{\mathcal{H}}_N - \partial_2 (\widehat{\mathcal{H}}_2 \varphi) - \partial_3 (\widehat{\mathcal{H}}_3 \varphi) \end{pmatrix}_{|x_1=0} = g, \end{aligned} \quad (2.10)$$

where $\widehat{v}_N = \widehat{v}_1 - \widehat{v}_2 \partial_2 \widehat{\Psi} - \widehat{v}_3 \partial_3 \widehat{\Psi}$, $\dot{\mathcal{H}}_N = \dot{\mathcal{H}}_1 - \dot{\mathcal{H}}_2 \partial_2 \widehat{\Psi} - \dot{\mathcal{H}}_3 \partial_3 \widehat{\Psi}$, and

$$[\partial_1 \dot{q}] = (\partial_1 \dot{q})|_{x_1=0} - (\dot{\mathcal{H}} \cdot \partial_1 \dot{\mathcal{H}})|_{x_1=0}.$$

We used assumption (2.5) taken at $x_1 = 0$ while writing down the last boundary condition in (2.10).

As in [1, 5, 21], we drop the zeroth-order terms in Ψ in (2.9) and consider the effective linear operators

$$\mathbb{P}'_e(\widehat{U}, \widehat{\Psi})\dot{U} := P(\widehat{U}, \widehat{\Psi})\dot{U} + \mathcal{C}(\widehat{U}, \widehat{\Psi})\dot{U} = f.$$

In the future nonlinear analysis the dropped terms in (2.9) should be considered as error terms. The new form of our linearized problem for $(\dot{U}, \dot{\mathcal{H}}, \varphi)$ reads:

$$\widehat{A}_0 \partial_t \dot{U} + \sum_{j=1}^3 \widehat{A}_j \partial_j \dot{U} + \widehat{\mathcal{C}} \dot{U} = f \quad \text{in } Q_T^+, \quad (2.11a)$$

$$\nabla \times \dot{\mathcal{H}} = \chi, \quad \operatorname{div} \dot{\mathcal{H}} = \Xi \quad \text{in } Q_T^-, \quad (2.11b)$$

$$\partial_t \varphi = \widehat{v}_N - \widehat{v}_2 \partial_2 \varphi - \widehat{v}_3 \partial_3 \varphi + \varphi \partial_1 \widehat{v}_N + g_1, \quad (2.11c)$$

$$\dot{q} = \dot{\mathcal{H}} \cdot \dot{\mathcal{H}} - [\partial_1 \dot{q}] \varphi + g_2, \quad (2.11d)$$

$$\dot{\mathcal{H}}_N = \partial_2 (\widehat{\mathcal{H}}_2 \varphi) + \partial_3 (\widehat{\mathcal{H}}_3 \varphi) + g_3 \quad \text{on } \omega_T, \quad (2.11e)$$

$$(\dot{U}, \dot{\mathcal{H}}, \varphi) = 0 \quad \text{for } t < 0, \quad (2.11f)$$

where

$$\widehat{A}_\alpha =: A_\alpha(\widehat{U}), \quad \alpha = 0, 2, 3, \quad \widehat{A}_1 =: \widetilde{A}_1(\widehat{U}, \widehat{\Psi}), \quad \widehat{\mathcal{C}} := \mathcal{C}(\widehat{U}, \widehat{\Psi}),$$

$$\dot{\mathcal{H}} = (\dot{\mathcal{H}}_1 \partial_1 \widehat{\Phi}_1, \dot{\mathcal{H}}_{\tau_2}, \dot{\mathcal{H}}_{\tau_3}), \quad \dot{\mathcal{H}} = (\dot{\mathcal{H}}_N, \dot{\mathcal{H}}_2 \partial_1 \widehat{\Phi}_1, \dot{\mathcal{H}}_3 \partial_1 \widehat{\Phi}_1),$$

$$\dot{\mathcal{H}}_N = \dot{\mathcal{H}}_1 - \dot{\mathcal{H}}_2 \partial_2 \widehat{\Psi} - \dot{\mathcal{H}}_3 \partial_3 \widehat{\Psi}, \quad \dot{\mathcal{H}}_{\tau_i} = \dot{\mathcal{H}}_1 \partial_i \widehat{\Psi} + \dot{\mathcal{H}}_i, \quad i = 2, 3.$$

The source term χ of the first equation in (2.11b) should satisfy the constraint $\operatorname{div} \chi = 0$. For the resolution of the elliptic problem (2.11b), (2.11e) the data Ξ, g_3 must satisfy the necessary compatibility condition

$$\int_{\Omega^-} \Xi \, dx = \int_{\Gamma} g_3 \, dx', \quad (2.12)$$

which follows from the double integration by parts

$$\int_{\Omega^-} \mathcal{E} \, dx = \int_{\Omega^-} \operatorname{div} \dot{\mathfrak{h}} \, dx = \int_{\Gamma} \dot{\mathfrak{h}}_1 \, dx' = \int_{\Gamma} \{\partial_2(\widehat{\mathcal{H}}_2\varphi) + \partial_3(\widehat{\mathcal{H}}_3\varphi) + g_3\} \, dx' = \int_{\Gamma} g_3 \, dx'.$$

We assume that the source terms f, χ, \mathcal{E} and the boundary datum g vanish in the past and consider the case of zero initial data. We postpone the case of nonzero initial data to the nonlinear analysis (see e.g. [5, 21]).

2.3 Reduction to homogeneous constraints in the “vacuum part”

We decompose $\dot{\mathcal{H}}$ in (2.11) as $\dot{\mathcal{H}} = \dot{\mathcal{H}}' + \dot{\mathcal{H}}''$ (and accordingly $\dot{\mathfrak{H}} = \dot{\mathfrak{H}}' + \dot{\mathfrak{H}}''$, $\dot{\mathfrak{h}} = \dot{\mathfrak{h}}' + \dot{\mathfrak{h}}''$), where \mathcal{H}'' is required to solve for each t the elliptic problem

$$\begin{aligned} \nabla \times \mathfrak{H}'' &= \chi, & \operatorname{div} \mathfrak{h}'' &= \mathcal{E} & \text{in } \Omega^-, \\ \mathfrak{h}_1'' &= \mathcal{H}_N'' = g_3 & & & \text{on } \Gamma. \end{aligned} \quad (2.13)$$

The source term χ of the first equation should satisfy the constraint $\operatorname{div} \chi = 0$. For the resolution of (2.13) the data \mathcal{E}, g_3 must satisfy the necessary compatibility condition (2.12). By classical results of the elliptic theory we have the following result.

LEMMA 2.1 Assume that the data (χ, \mathcal{E}, g_3) in (2.13), vanishing in appropriate way as x goes to infinity, satisfy the constraint $\operatorname{div} \chi = 0$ and the compatibility condition (2.12). Then there exists a unique solution \mathcal{H}'' of (2.13) vanishing at infinity.

REMARK In the statement of the lemma above we intentionally leave unspecified the description of the regularity and the behavior at infinity of the data and consequently of the solution. This point will be faced in the forthcoming paper on the resolution of the nonlinear problem.

Given \mathcal{H}'' , now we look for \mathcal{H}' such that

$$\begin{aligned} \nabla \times \mathfrak{H}' &= 0, & \operatorname{div} \mathfrak{h}' &= 0 & \text{in } Q_T^-, \\ q &= \widehat{\mathcal{H}} \cdot \mathcal{H}' - [\partial_1 \hat{q}] \varphi + g_2', & & & \\ \mathcal{H}'_N &= \partial_2(\widehat{\mathcal{H}}_2\varphi) + \partial_3(\widehat{\mathcal{H}}_3\varphi) & & & \text{on } \omega_T, \end{aligned} \quad (2.14)$$

where we have denoted $g_2' = g_2 + \widehat{\mathcal{H}} \cdot \mathcal{H}''$. If \mathcal{H}'' solves (2.13) and \mathcal{H}' is a solution of (2.14) then $\dot{\mathcal{H}} = \dot{\mathcal{H}}' + \dot{\mathcal{H}}''$ clearly solves (2.11b), (2.11d), (2.11e).

From (2.11), (2.14), the new form of the reduced linearized problem with unknowns (U, \mathcal{H}') reads (we drop for convenience the $'$ in \mathcal{H}', g_2')

$$\widehat{A}_0 \partial_t \dot{U} + \sum_{j=1}^3 \widehat{A}_j \partial_j \dot{U} + \widehat{\mathcal{C}} \dot{U} = f \quad \text{in } Q_T^+, \quad (2.15a)$$

$$\nabla \times \mathfrak{H} = 0, \quad \operatorname{div} \mathfrak{h} = 0 \quad \text{in } Q_T^-, \quad (2.15b)$$

$$\partial_t \varphi = \dot{v}_N - \dot{v}_2 \partial_2 \varphi - \dot{v}_3 \partial_3 \varphi + \varphi \partial_1 \dot{v}_N + g_1, \quad (2.15c)$$

$$\dot{q} = \widehat{\mathcal{H}} \cdot \mathcal{H} - [\partial_1 \hat{q}] \varphi + g_2, \quad (2.15d)$$

$$\mathcal{H}_N = \partial_2(\widehat{\mathcal{H}}_2\varphi) + \partial_3(\widehat{\mathcal{H}}_3\varphi) \quad \text{on } \omega_T, \quad (2.15e)$$

$$(\dot{U}, \mathcal{H}, \varphi) = 0 \quad \text{for } t < 0. \quad (2.15f)$$

2.4 Reduction to homogeneous constraints in the “plasma part”

From problem (2.15) we can deduce nonhomogeneous equations associated with the divergence constraint $\operatorname{div} \dot{h} = 0$ and the “redundant” boundary conditions $\dot{H}_N|_{x_1=0} = 0$ for the nonlinear problem. More precisely, with reference to [21, Proposition 2] for the proof, we have the following.

PROPOSITION 2.2 ([21]) Let the basic state (2.1) satisfies assumptions (2.2)–(2.7). Then solutions of problem (2.15) satisfy

$$\operatorname{div} \dot{h} = r \quad \text{in } Q_T^+, \quad (2.16)$$

$$\widehat{H}_2 \partial_2 \varphi + \widehat{H}_3 \partial_3 \varphi - \dot{H}_N - \varphi \partial_1 \widehat{H}_N = R \quad \text{on } \omega_T. \quad (2.17)$$

Here

$$\dot{h} = (\dot{H}_N, \dot{H}_2 \partial_1 \widehat{\Phi}_1, \dot{H}_3 \partial_1 \widehat{\Phi}_1), \quad \dot{H}_N = \dot{H}_1 - \dot{H}_2 \partial_2 \widehat{\Psi} - \dot{H}_3 \partial_3 \widehat{\Psi}.$$

The functions $r = r(t, x)$ and $R = R(t, x')$, which vanish in the past, are determined by the source terms and the basic state as solutions to the linear inhomogeneous equations

$$\partial_t a + \frac{1}{\partial_1 \widehat{\Phi}_1} \{ \hat{w} \cdot \nabla a + a \operatorname{div} \hat{w} \} = \mathfrak{F}_H \quad \text{in } Q_T^+, \quad (2.18)$$

$$\partial_t R + \hat{v}_2 \partial_2 R + \hat{v}_3 \partial_3 R + (\partial_2 \hat{v}_2 + \partial_3 \hat{v}_3) R = \mathcal{Q} \quad \text{on } \omega_T, \quad (2.19)$$

where $a = r / \partial_1 \widehat{\Phi}_1$, $\mathfrak{F}_H = (\operatorname{div} f_H) / \partial_1 \widehat{\Phi}_1$,

$$f_H = (f_N, f_6, f_7), \quad f_N = f_5 - f_6 \partial_2 \widehat{\Psi} - f_7 \partial_3 \widehat{\Psi}, \quad \mathcal{Q} = \{ \partial_2 (\widehat{H}_2 g_1) + \partial_3 (\widehat{H}_3 g_1) - f_N \} |_{x_1=0}.$$

Let us reduce (2.15) to a problem with homogeneous boundary conditions (2.15c), (2.15d) (i.e., $g_1 = g_2 = 0$) and homogeneous constraints (2.16) and (2.17) (i.e., $r = R = 0$). More precisely, we describe a “lifting” function as follows:

$$\widetilde{U} = (\widetilde{q}, \widetilde{v}_1, 0, 0, \widetilde{H}, 0),$$

where $\widetilde{q} = g_2$, $\widetilde{v}_1 = -g_1$ on ω_T , and where \widetilde{H} solves the equation for \dot{H} contained in (2.15a) with $\dot{v} = 0$:

$$\partial_t \widetilde{H} + \frac{1}{\partial_1 \widehat{\Phi}_1} \{ (\hat{w} \cdot \nabla) \widetilde{H} - (\widetilde{h} \cdot \nabla) \hat{v} + \widetilde{H} \operatorname{div} \hat{w} \} = f_H \quad \text{in } Q_T^+, \quad (2.20)$$

where $\widetilde{h} = (\widetilde{H}_1 - \widetilde{H}_2 \partial_2 \widehat{\Psi} - \widetilde{H}_3 \partial_3 \widehat{\Psi}, \widetilde{H}_2, \widetilde{H}_3)$, $f_H = (f_5, f_6, f_7)$. It is very important that, in view of (2.6), we have $\hat{w}_1|_{x_1=0} = 0$; therefore the linear equation (2.20) does not need any boundary condition. Then the new unknown

$$U^\natural = \dot{U} - \widetilde{U}, \quad \mathcal{H}^\natural = \mathcal{H} \quad (2.21)$$

satisfies problem (2.15) with $f = F$, where

$$F = (F_1, \dots, F_8) = f - \mathbb{P}'_e(\widehat{U}, \widehat{\Psi}) \widetilde{U}.$$

In view of (2.20), $F_H = (F_5, F_6, F_7) = 0$, and it follows from Proposition 2.2 that U^\natural satisfies (2.16) and (2.17) with $r = R = 0$.

Dropping for convenience the indices \mathfrak{h} in (2.21), the new form of our reduced linearized problem now reads

$$\widehat{A}_0 \partial_t U + \sum_{j=1}^3 \widehat{A}_j \partial_j U + \widehat{\mathcal{C}} U = F \quad \text{in } Q_T^+, \quad (2.22a)$$

$$\nabla \times \mathfrak{H} = 0, \quad \operatorname{div} \mathfrak{h} = 0 \quad \text{in } Q_T^-, \quad (2.22b)$$

$$\partial_t \varphi = v_N - \widehat{v}_2 \partial_2 \varphi - \widehat{v}_3 \partial_3 \varphi + \varphi \partial_1 \widehat{v}_N, \quad (2.22c)$$

$$q = \widehat{\mathcal{K}} \cdot \mathcal{K} - [\partial_1 \widehat{q}] \varphi, \quad (2.22d)$$

$$\mathcal{K}_N = \partial_2 (\widehat{\mathcal{K}}_2 \varphi) + \partial_3 (\widehat{\mathcal{K}}_3 \varphi) \quad \text{on } \omega_T, \quad (2.22e)$$

$$(U, \mathcal{K}, \varphi) = 0 \quad \text{for } t < 0. \quad (2.22f)$$

and solutions should satisfy

$$\operatorname{div} h = 0 \quad \text{in } Q_T^+, \quad (2.23)$$

$$H_N = \widehat{H}_2 \partial_2 \varphi + \widehat{H}_3 \partial_3 \varphi - \varphi \partial_1 \widehat{H}_N \quad \text{on } \omega_T. \quad (2.24)$$

257 All the notations here for U and \mathcal{K} (e.g., h , \mathfrak{H} , \mathfrak{h} , etc.) are analogous to the corresponding ones for
258 \dot{U} and $\dot{\mathcal{K}}$ introduced above.

259 2.5 An equivalent formulation of (2.22)

In the following analysis it is convenient to make use of different “plasma” variables and an equivalent form of equations (2.22a). We define the matrix

$$\widehat{\eta} = \begin{pmatrix} 1 & -\partial_2 \widehat{\Psi} & -\partial_3 \widehat{\Psi} \\ 0 & \partial_1 \widehat{\Phi}_1 & 0 \\ 0 & 0 & \partial_1 \widehat{\Phi}_1 \end{pmatrix}.$$

260 It follows that

$$261 \quad u = (v_N, v_2 \partial_1 \widehat{\Phi}_1, v_3 \partial_1 \widehat{\Phi}_1) = \widehat{\eta} v, \quad h = (H_N, H_2 \partial_1 \widehat{\Phi}_1, H_3 \partial_1 \widehat{\Phi}_1) = \widehat{\eta} H. \quad (2.25)$$

Multiplying (2.22a) on the left side by the matrix

$$\widehat{R} = \begin{pmatrix} 1 & \underline{0} & \underline{0} & 0 \\ \underline{0}^T & \widehat{\eta} & 0_3 & \underline{0}^T \\ \underline{0}^T & 0_3 & \widehat{\eta} & \underline{0}^T \\ 0 & \underline{0}^T & \underline{0}^T & 1 \end{pmatrix},$$

after some calculations we get the symmetric hyperbolic system for the new vector of unknowns $\mathcal{U} = (q, u, h, S)$ (compare with (1.6), (2.22a)):

$$\begin{aligned} & \partial_1 \widehat{\Phi}_1 \begin{pmatrix} \hat{\rho}_p/\hat{\rho} & \underline{0} & -(\hat{\rho}_p/\hat{\rho})\hat{h} & 0 \\ \underline{0}^T & \hat{\rho}\hat{a}_0 & 0_3 & \underline{0}^T \\ -(\hat{\rho}_p/\hat{\rho})\hat{h}^T & 0_3 & \hat{a}_0 + (\hat{\rho}_p/\hat{\rho})\hat{h} \otimes \hat{h} & \underline{0}^T \\ 0 & \underline{0} & \underline{0} & 1 \end{pmatrix} \partial_t \begin{pmatrix} q \\ u \\ h \\ S \end{pmatrix} + \begin{pmatrix} 0 & \nabla \cdot & \underline{0} & 0 \\ \nabla & 0_3 & 0_3 & \underline{0}^T \\ \underline{0}^T & 0_3 & 0_3 & \underline{0}^T \\ 0 & \underline{0} & \underline{0} & 0 \end{pmatrix} \begin{pmatrix} q \\ u \\ h \\ S \end{pmatrix} \\ & + \partial_1 \widehat{\Phi}_1 \begin{pmatrix} (\hat{\rho}_p/\hat{\rho})\hat{w} \cdot \nabla & \nabla \cdot & -(\hat{\rho}_p/\hat{\rho})\hat{h}\hat{w} \cdot \nabla & 0 \\ \nabla & \hat{\rho}\hat{a}_0\hat{w} \cdot \nabla & -\hat{a}_0\hat{h} \cdot \nabla & \underline{0}^T \\ -(\hat{\rho}_p/\hat{\rho})\hat{h}^T \hat{w} \cdot \nabla & -\hat{a}_0\hat{h} \cdot \nabla & (\hat{a}_0 + (\hat{\rho}_p/\hat{\rho})\hat{h} \otimes \hat{h})\hat{w} \cdot \nabla & \underline{0}^T \\ 0 & \underline{0} & \underline{0} & \hat{w} \cdot \nabla \end{pmatrix} \begin{pmatrix} q \\ u \\ h \\ S \end{pmatrix} + \widehat{\mathcal{C}}' \mathcal{U} = \mathcal{F}, \end{aligned} \quad (2.26)$$

where \hat{a}_0 is the symmetric and positive definite matrix

$$\hat{a}_0 = (\hat{\eta}^{-1})^T \hat{\eta}^{-1},$$

with a new matrix $\widehat{\mathcal{C}}'$ in the zero-order term (whose precise form has no importance) and where we have set $\mathcal{F} = \partial_1 \widehat{\Phi}_1 \widehat{R}F$. We write system (2.26) in compact form as

$$\widehat{\mathcal{G}}_0 \partial_t \mathcal{U} + \sum_{j=1}^3 (\widehat{\mathcal{G}}_j + \mathcal{E}_{1j+1}) \partial_j \mathcal{U} + \widehat{\mathcal{C}}' \mathcal{U} = \mathcal{F}, \quad (2.27)$$

where

$$\begin{aligned} \mathcal{E}_{12} &= \begin{pmatrix} 0 & 1 & 0 & 0 & \dots & 0 \\ 1 & 0 & 0 & 0 & \dots & 0 \\ 0 & 0 & 0 & 0 & \dots & 0 \\ 0 & 0 & 0 & 0 & \dots & 0 \\ \vdots & \vdots & \vdots & \vdots & & \vdots \\ 0 & 0 & 0 & 0 & \dots & 0 \end{pmatrix}, \quad \mathcal{E}_{13} = \begin{pmatrix} 0 & 0 & 1 & 0 & \dots & 0 \\ 0 & 0 & 0 & 0 & \dots & 0 \\ 1 & 0 & 0 & 0 & \dots & 0 \\ 0 & 0 & 0 & 0 & \dots & 0 \\ \vdots & \vdots & \vdots & \vdots & & \vdots \\ 0 & 0 & 0 & 0 & \dots & 0 \end{pmatrix}, \\ \mathcal{E}_{14} &= \begin{pmatrix} 0 & 0 & 0 & 1 & \dots & 0 \\ 0 & 0 & 0 & 0 & \dots & 0 \\ 0 & 0 & 0 & 0 & \dots & 0 \\ 1 & 0 & 0 & 0 & \dots & 0 \\ \vdots & \vdots & \vdots & \vdots & & \vdots \\ 0 & 0 & 0 & 0 & \dots & 0 \end{pmatrix}. \end{aligned}$$

The formulation (2.27) has the advantage of the form of the boundary matrix of the system $\widehat{\mathcal{G}}_1 + \mathcal{E}_{12}$, with

$$\widehat{\mathcal{G}}_1 = 0 \quad \text{on } \omega_T, \quad (2.28)$$

because $\hat{w}_1 = \hat{h}_1 = 0$, and \mathcal{E}_{12} a constant matrix. Thus system (2.27) is symmetric hyperbolic with characteristic boundary of constant multiplicity (see [18–20] for maximally dissipative boundary

conditions). Thus, the final form of our reduced linearized problem is

$$\widehat{\mathcal{Q}}_0 \partial_t \mathcal{U} + \sum_{j=1}^3 (\widehat{\mathcal{Q}}_j + \varepsilon_{1j+1}) \partial_j \mathcal{U} + \widehat{\mathcal{C}}' \mathcal{U} = \mathcal{F}, \quad \text{in } Q_T^+, \quad (2.29a)$$

$$\nabla \times \mathfrak{H} = 0, \quad \operatorname{div} \mathfrak{h} = 0 \quad \text{in } Q_T^-, \quad (2.29b)$$

$$\partial_t \varphi = u_1 - \hat{v}_2 \partial_2 \varphi - \hat{v}_3 \partial_3 \varphi + \varphi \partial_1 \hat{v}_N, \quad (2.29c)$$

$$q = \widehat{\mathcal{H}} \cdot \mathcal{H} - [\partial_1 \hat{q}] \varphi, \quad (2.29d)$$

$$\mathcal{H}_N = \partial_2 (\widehat{\mathcal{H}}_2 \varphi) + \partial_3 (\widehat{\mathcal{H}}_3 \varphi) \quad \text{on } \omega_T, \quad (2.29e)$$

$$(\mathcal{U}, \mathcal{H}, \varphi) = 0 \quad \text{for } t < 0, \quad (2.29f)$$

under the constraints (2.23), (2.24).

3. Function spaces

Now we introduce the main function spaces to be used in the following. Let us denote

$$Q^\pm := \mathbb{R}_t \times \Omega^\pm, \quad \omega := \mathbb{R}_t \times \Gamma. \quad (3.1)$$

3.1 Weighted Sobolev spaces

For $\gamma \geq 1$ and $s \in \mathbb{R}$, we set

$$\lambda^{s,\gamma}(\xi) := (\gamma^2 + |\xi|^2)^{s/2}$$

and, in particular, $\lambda^{s,1} := \lambda^s$.

Throughout the paper, for real $\gamma \geq 1$ and $n \geq 2$, $H_\gamma^s(\mathbb{R}^n)$ will denote the Sobolev space of order s , equipped with the γ -dependent norm $\|\cdot\|_{s,\gamma}$ defined by

$$\|u\|_{s,\gamma}^2 := (2\pi)^{-n} \int_{\mathbb{R}^n} \lambda^{2s,\gamma}(\xi) |\widehat{u}(\xi)|^2 d\xi, \quad (3.2)$$

\widehat{u} being the Fourier transform of u . The norms defined by (3.2), with different values of the parameter γ , are equivalent each other. For $\gamma = 1$ we set for brevity $\|\cdot\|_s := \|\cdot\|_{s,1}$ (and, accordingly, the standard Sobolev space $H^s(\mathbb{R}^n) := H_1^s(\mathbb{R}^n)$). For $s \in \mathbb{N}$, the norm in (3.2) turns to be equivalent, *uniformly with respect to* γ , to the norm $\|\cdot\|_{H_\gamma^s(\mathbb{R}^n)}$ defined by

$$\|u\|_{H_\gamma^s(\mathbb{R}^n)}^2 := \sum_{|\alpha| \leq s} \gamma^{2(s-|\alpha|)} \|\partial^\alpha u\|_{L^2(\mathbb{R}^n)}^2.$$

For functions defined over Q_T^- we will consider the weighted Sobolev spaces $H_\gamma^m(Q_T^-)$ equipped with the γ -dependent norm

$$\|u\|_{H_\gamma^m(Q_T^-)}^2 := \sum_{|\alpha| \leq m} \gamma^{2(m-|\alpha|)} \|\partial^\alpha u\|_{L^2(Q_T^-)}^2.$$

Similar weighted Sobolev spaces will be considered for functions defined on Q^- .

3.2 Conormal Sobolev spaces

Let us introduce some classes of function spaces of Sobolev type, defined over the half-space Q_T^+ . For $j = 0, \dots, 3$, we set

$$Z_0 = \partial_t, \quad Z_1 := \sigma(x_1)\partial_1, \quad Z_j := \partial_j, \quad \text{for } j = 2, 3,$$

where $\sigma(x_1) \in C^\infty(\mathbb{R}_+)$ is a monotone increasing function such that $\sigma(x_1) = x_1$ in a neighborhood of the origin and $\sigma(x_1) = 1$ for x_1 large enough. Then, for every multi-index $\alpha = (\alpha_0, \dots, \alpha_3) \in \mathbb{N}^4$, the *conormal* derivative Z^α is defined by

$$Z^\alpha := Z_0^{\alpha_0} \dots Z_3^{\alpha_3};$$

we also write $\partial^\alpha = \partial_0^{\alpha_0} \dots \partial_3^{\alpha_3}$ for the usual partial derivative corresponding to α .

Given an integer $m \geq 1$, the *conormal Sobolev space* $H_{tan}^m(Q_T^+)$ is defined as the set of functions $u \in L^2(Q_T^+)$ such that $Z^\alpha u \in L^2(Q_T^+)$, for all multi-indices α with $|\alpha| \leq m$ (see [15, 16]). Agreeing with the notations set for the usual Sobolev spaces, for $\gamma \geq 1$, $H_{tan,\gamma}^m(Q_T^+)$ will denote the conormal space of order m equipped with the γ -depending norm

$$\|u\|_{H_{tan,\gamma}^m(Q_T^+)}^2 := \sum_{|\alpha| \leq m} \gamma^{2(m-|\alpha|)} \|Z^\alpha u\|_{L^2(Q_T^+)}^2 \quad (3.3)$$

and we have $H_{tan}^m(Q_T^+) := H_{tan,1}^m(Q_T^+)$. Similar conormal Sobolev spaces with γ -depending norms will be considered for functions defined on Q^+ .

We will use the same notation for spaces of scalar and vector-valued functions.

4. The main result

We are now in a position to state the main result of this paper. Recall that $\mathcal{U} = (q, u, h, S)$, where u and h were defined in (2.25).

MAIN THEOREM 4.1 Let $T > 0$. Let the basic state (2.1) satisfies assumptions (2.2)–(2.7) and

$$|\widehat{H} \times \widehat{\mathcal{H}}| \geq \delta > 0 \quad \text{on } \omega_T, \quad (4.1)$$

where δ is a fixed constant. There exists $\gamma_0 \geq 1$ such that for all $\gamma \geq \gamma_0$ and for all $\mathfrak{F}_\gamma \in H_{tan,\gamma}^1(Q_T^+)$, vanishing in the past, namely for $t < 0$, problem (2.29) has a unique solution $(\mathcal{U}_\gamma, \mathcal{H}_\gamma, \varphi_\gamma) \in H_{tan,\gamma}^1(Q_T^+) \times H_\gamma^1(Q_T^-) \times H_\gamma^1(\omega_T)$ with trace $(q_\gamma, u_{1\gamma}, h_{1\gamma})|_{\omega_T} \in H_\gamma^{1/2}(\omega_T)$, $\mathcal{H}_\gamma|_{\omega_T} \in H_\gamma^{1/2}(\omega_T)$. Moreover, the solution obeys the a priori estimate

$$\begin{aligned} \gamma \left(\|\mathcal{U}_\gamma\|_{H_{tan,\gamma}^1(Q_T^+)}^2 + \|\mathcal{H}_\gamma\|_{H_\gamma^1(Q_T^-)}^2 + \|(q_\gamma, u_{1\gamma}, h_{1\gamma})|_{\omega_T}\|_{H_\gamma^{1/2}(\omega_T)}^2 + \|\mathcal{H}_\gamma|_{\omega_T}\|_{H_\gamma^{1/2}(\omega_T)}^2 \right) \\ + \gamma^2 \|\varphi_\gamma\|_{H_\gamma^1(\omega_T)}^2 \leq \frac{C}{\gamma} \|\mathfrak{F}_\gamma\|_{H_{tan,\gamma}^1(Q_T^+)}^2, \end{aligned} \quad (4.2)$$

where we have set $\mathcal{U}_\gamma = e^{-\gamma t} \mathcal{U}$, $\mathcal{H}_\gamma = e^{-\gamma t} \mathcal{H}$, $\varphi_\gamma = e^{-\gamma t} \varphi$ and so on. Here $C = C(K, T, \delta) > 0$ is a constant independent of the data \mathfrak{F} and γ .

The a priori estimate (4.2) improves the similar estimate firstly proved in [23].

REMARK Strictly speaking, the uniqueness of the solution to problem (2.29) follows from the a priori estimate (42) derived in [23], provided that our solution belongs to H^2 . We do not present here a formal proof of the existence of solutions with a higher degree of regularity (in particular, H^2) and postpone this part to the future work on the nonlinear problem (see e.g. [5, 21]).

The remainder of the paper is organized as follows. In Section 5 we introduce a fully hyperbolic regularization of the coupled hyperbolic-elliptic system (2.29). In Section 6 we show an a priori estimate of solutions uniform in the small parameter ε of regularization. In Section 7 we show the well-posedness of the hyperbolic regularization and in Section 8 we conclude the proof of Theorem 4.1 by passing to the limit as $\varepsilon \rightarrow 0$. Sections 9, 10, and 11 are devoted to the proof of some technical results.

5. Hyperbolic regularization of the reduced problem

The problem (2.29) is a nonstandard initial-boundary value problem for a coupled hyperbolic-elliptic system. For its resolution we introduce a “hyperbolic” regularization of the elliptic system (2.29b). We will prove the existence of solutions for such regularized problem by referring to the well-posedness theory for linear symmetric hyperbolic systems with characteristic boundary and maximally nonnegative boundary conditions [19, 20]. After showing suitable a priori estimate uniform in ε , we will pass to the limit as $\varepsilon \rightarrow 0$, to get the solution of (2.29).

The regularization of problem (2.29) is inspired by a corresponding problem in relativistic MHD [24]. In our non-relativistic case the displacement current $(1/c)\partial_t E$ is neglected in the vacuum Maxwell equations, where c is the speed of light and E is the electric field. Now, in some sense, we restore this neglected term. Namely, we consider a “hyperbolic” regularization of the elliptic system (2.29b) by introducing a new auxiliary unknown E^ε which plays a role of the vacuum electric field, and the small parameter of regularization ε is associated with the physical parameter $1/c$. We also regularize the second boundary condition in (2.29d) and introduce two boundary conditions for the unknown E^ε .

Let us denote $V^\varepsilon = (\mathcal{H}^\varepsilon, E^\varepsilon)$. Given a small parameter $\varepsilon > 0$, we consider the following regularized problem for the unknown $(\mathcal{U}^\varepsilon, V^\varepsilon, \varphi^\varepsilon)$:

$$\widehat{\alpha}_0 \partial_t \mathcal{U}^\varepsilon + \sum_{j=1}^3 (\widehat{\alpha}_j + \varepsilon_{1j+1}) \partial_j \mathcal{U}^\varepsilon + \widehat{\mathcal{C}}' \mathcal{U}^\varepsilon = \mathfrak{F} \quad \text{in } Q_T^+, \quad (5.1a)$$

$$\varepsilon \partial_t \mathfrak{h}^\varepsilon + \nabla \times \mathfrak{E}^\varepsilon = 0, \quad \varepsilon \partial_t \mathfrak{e}^\varepsilon - \nabla \times \mathfrak{H}^\varepsilon = 0 \quad \text{in } Q_T^-, \quad (5.1b)$$

$$\partial_t \varphi^\varepsilon = u_1^\varepsilon - \widehat{v}_2 \partial_2 \varphi^\varepsilon - \widehat{v}_3 \partial_3 \varphi^\varepsilon + \varphi^\varepsilon \partial_1 \widehat{v}_N, \quad (5.1c)$$

$$q^\varepsilon = \widehat{\mathcal{H}} \cdot \mathcal{H}^\varepsilon - [\partial_1 \widehat{q}] \varphi^\varepsilon - \varepsilon \widehat{E} \cdot E^\varepsilon, \quad (5.1d)$$

$$E_{\tau_2}^\varepsilon = \varepsilon \partial_t (\widehat{\mathcal{H}}_3 \varphi^\varepsilon) - \varepsilon \partial_2 (\widehat{E}_1 \varphi^\varepsilon), \quad (5.1e)$$

$$E_{\tau_3}^\varepsilon = -\varepsilon \partial_t (\widehat{\mathcal{H}}_2 \varphi^\varepsilon) - \varepsilon \partial_3 (\widehat{E}_1 \varphi^\varepsilon) \quad \text{on } \omega_T, \quad (5.1f)$$

$$(\mathcal{U}^\varepsilon, V^\varepsilon, \varphi^\varepsilon) = 0 \quad \text{for } t < 0, \quad (5.1g)$$

where

$$\begin{aligned} E^\varepsilon &= (E_1^\varepsilon, E_2^\varepsilon, E_3^\varepsilon), \quad \widehat{E} = (\widehat{E}_1, \widehat{E}_2, \widehat{E}_3), \quad \mathfrak{E}^\varepsilon = (E_1^\varepsilon \partial_1 \widehat{\Phi}_1, E_{\tau_2}^\varepsilon, E_{\tau_3}^\varepsilon), \\ \mathfrak{e}^\varepsilon &= (E_N^\varepsilon, E_2^\varepsilon \partial_1 \widehat{\Phi}_1, E_3^\varepsilon \partial_1 \widehat{\Phi}_1), \quad E_N^\varepsilon = E_1^\varepsilon - E_2^\varepsilon \partial_2 \widehat{\Psi} - E_3^\varepsilon \partial_3 \widehat{\Psi}, \\ E_{\tau_k}^\varepsilon &= E_1^\varepsilon \partial_k \widehat{\Psi} + E_k^\varepsilon, \quad k = 2, 3, \end{aligned}$$

the coefficients \widehat{E}_j are given functions which will be chosen later on. All the other notations for \mathcal{H}^ε (e.g., $\mathfrak{H}^\varepsilon, \mathfrak{h}^\varepsilon$) are analogous to those for \mathcal{H} .

If $\Psi = 0, \Phi_1 = x_1$, then $\mathfrak{h}^\varepsilon = \mathfrak{H}^\varepsilon = \mathcal{H}^\varepsilon, \mathfrak{e}^\varepsilon = \mathfrak{E}^\varepsilon = E^\varepsilon$, and when $\varepsilon = 1$ (5.1b) turns out to be nothing else than the Maxwell equations.

It is noteworthy that solutions to problem (5.1) satisfy

$$\operatorname{div} h^\varepsilon = 0 \quad \text{in } Q_T^+, \quad (5.2)$$

$$\operatorname{div} \mathfrak{h}^\varepsilon = 0, \quad \operatorname{div} \mathfrak{e}^\varepsilon = 0 \quad \text{in } Q_T^-, \quad (5.3)$$

$$h_1^\varepsilon = \widehat{H}_2 \partial_2 \varphi^\varepsilon + \widehat{H}_3 \partial_3 \varphi^\varepsilon - \varphi^\varepsilon \partial_1 \widehat{H}_N, \quad (5.4)$$

$$\mathcal{H}_N^\varepsilon = \partial_2 (\widehat{\mathcal{H}}_2 \varphi^\varepsilon) + \partial_3 (\widehat{\mathcal{H}}_3 \varphi^\varepsilon) \quad \text{on } \omega_T, \quad (5.5)$$

because (5.2)–(5.5) are just restrictions on the initial data which are automatically satisfied in view of (5.1g). Indeed, the derivation of (5.2) and (5.4) is absolutely the same as that of (2.23) and (2.24). Equations (5.3) trivially follow from (5.1b), (5.1g). Moreover, condition (5.5) is obtained by considering the first component of the first equation in (5.1b) at $x_1 = 0$ and taking into account (5.1e)–(5.1g).

5.1 An equivalent formulation of (5.1)

In the following analysis it is convenient to make use of a different formulation of the approximating problem (5.1), as far as the vacuum part is concerned.

First we introduce the matrices which are coefficients of the space derivatives in (5.1b) (for $\varepsilon = 1$ the matrices below are those for the vacuum Maxwell equations):

$$B_1^\varepsilon = \varepsilon^{-1} \begin{pmatrix} 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & -1 \\ 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 & 0 & 0 \end{pmatrix}, \quad B_2^\varepsilon = \varepsilon^{-1} \begin{pmatrix} 0 & 0 & 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 & 0 & 0 \end{pmatrix},$$

$$B_3^\varepsilon = \varepsilon^{-1} \begin{pmatrix} 0 & 0 & 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 \\ -1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \end{pmatrix}.$$

Then system (5.1b) can be written in terms of the “curved” unknown $W^\varepsilon = (\mathfrak{H}^\varepsilon, \mathfrak{E}^\varepsilon)$ as

$$B_0 \partial_t W^\varepsilon + \sum_{j=1}^3 B_j^\varepsilon \partial_j W^\varepsilon + B_4 W^\varepsilon = 0, \quad (5.6)$$

where

$$B_0 = (\partial_1 \widehat{\Phi}_1)^{-1} K K^\top > 0, \quad K = I_2 \otimes \hat{\eta}, \quad B_4 = \partial_t B_0,$$

and the matrices B_0 and K are found from the relations

$$\mathfrak{h}^\varepsilon = \hat{\eta} \mathcal{H}^\varepsilon = (\partial_1 \widehat{\Phi}_1)^{-1} \hat{\eta} \hat{\eta}^\top \mathfrak{H}^\varepsilon, \quad \mathfrak{e}^\varepsilon = \hat{\eta} E^\varepsilon = (\partial_1 \widehat{\Phi}_1)^{-1} \hat{\eta} \hat{\eta}^\top \mathfrak{E}^\varepsilon,$$

so that

$$\begin{pmatrix} \mathfrak{h}^\varepsilon \\ \mathfrak{e}^\varepsilon \end{pmatrix} = (\partial_1 \widehat{\Phi}_1)^{-1} \begin{pmatrix} \hat{\eta} \hat{\eta}^\top & 0_3 \\ 0_3 & \hat{\eta} \hat{\eta}^\top \end{pmatrix} \begin{pmatrix} \mathfrak{H}^\varepsilon \\ \mathfrak{E}^\varepsilon \end{pmatrix} = B_0 W^\varepsilon.$$

System (5.6) is symmetric hyperbolic. The convenience of the use of variables $(\mathfrak{H}^\varepsilon, \mathfrak{E}^\varepsilon)$ rather than $(\mathcal{H}^\varepsilon, E^\varepsilon)$ stays mainly in that the matrices B_j^ε of (5.6), containing the singular multiplier ε^{-1} , are constant.

Finally, we write the boundary conditions (5.1c)–(5.1f) in terms of $(\mathcal{U}^\varepsilon, W^\varepsilon)$, where we observe that (recalling that $\partial_1 \widehat{\Phi}_1 = 1$ on ω_T):

$$\begin{aligned} \widehat{\mathcal{H}} \cdot \mathcal{H}^\varepsilon &= \widehat{\mathcal{H}}_N \mathcal{H}_1^\varepsilon + \widehat{\mathcal{H}}_2 \mathcal{H}_{\tau_2}^\varepsilon + \widehat{\mathcal{H}}_3 \mathcal{H}_{\tau_3}^\varepsilon = \hat{\mathfrak{h}} \cdot \mathfrak{H}^\varepsilon, \\ \widehat{E} \cdot E^\varepsilon &= \widehat{E}_N E_1^\varepsilon + \widehat{E}_2 E_{\tau_2}^\varepsilon + \widehat{E}_3 E_{\tau_3}^\varepsilon = \hat{\mathfrak{e}} \cdot \mathfrak{E}^\varepsilon. \end{aligned} \quad (5.7)$$

Concerning the first line above in (5.7) we notice that $\hat{\mathfrak{h}}_1 = \widehat{\mathcal{H}}_N = 0$ on ω_T , so that $\mathfrak{H}_1^\varepsilon$ does not appear in the boundary condition.

From (5.6), (5.7) we get the new formulation of problem (5.1) for the unknowns $(\mathcal{U}^\varepsilon, W^\varepsilon)$:

$$\widehat{\mathcal{Q}}_0 \partial_t \mathcal{U}^\varepsilon + \sum_{j=1}^3 (\widehat{\mathcal{Q}}_j + \mathfrak{E}_{1j+1}) \partial_j \mathcal{U}^\varepsilon + \widehat{\mathcal{C}}' \mathcal{U}^\varepsilon = \mathfrak{F}, \quad \text{in } Q_T^+, \quad (5.8a)$$

$$B_0 \partial_t W^\varepsilon + \sum_{j=1}^3 B_j^\varepsilon \partial_j W^\varepsilon + B_4 W^\varepsilon = 0 \quad \text{in } Q_T^-, \quad (5.8b)$$

$$\partial_t \varphi^\varepsilon + \hat{v}_2 \partial_2 \varphi^\varepsilon + \hat{v}_3 \partial_3 \varphi^\varepsilon - \varphi^\varepsilon \partial_1 \hat{v}_N - u_1^\varepsilon = 0, \quad (5.8c)$$

$$q^\varepsilon + [\partial_1 \hat{q}] \varphi^\varepsilon - \hat{\mathfrak{h}} \cdot \mathfrak{H}^\varepsilon + \varepsilon \hat{\mathfrak{e}} \cdot \mathfrak{E}^\varepsilon = 0, \quad (5.8d)$$

$$\mathfrak{E}_2^\varepsilon - \varepsilon \partial_t (\widehat{\mathcal{H}}_3 \varphi^\varepsilon) + \varepsilon \partial_2 (\widehat{E}_1 \varphi^\varepsilon) = 0, \quad (5.8e)$$

$$\mathfrak{E}_3^\varepsilon + \varepsilon \partial_t (\widehat{\mathcal{H}}_2 \varphi^\varepsilon) + \varepsilon \partial_3 (\widehat{E}_1 \varphi^\varepsilon) = 0 \quad \text{on } \omega_T, \quad (5.8f)$$

$$(\mathcal{U}^\varepsilon, W^\varepsilon, \varphi^\varepsilon) = 0 \quad \text{for } t < 0. \quad (5.8g)$$

From (5.2)–(5.5) we get that solutions $(\mathcal{U}^\varepsilon, W^\varepsilon)$ to problem (5.8) satisfy

$$\operatorname{div} h^\varepsilon = 0 \quad \text{in } Q_T^+, \quad (5.9)$$

$$\operatorname{div} \mathfrak{h}^\varepsilon = 0, \quad \operatorname{div} \mathfrak{e}^\varepsilon = 0 \quad \text{in } Q_T^-, \quad (5.10)$$

$$h_1^\varepsilon = \widehat{H}_2 \partial_2 \varphi^\varepsilon + \widehat{H}_3 \partial_3 \varphi^\varepsilon - \varphi^\varepsilon \partial_1 \widehat{H}_N, \quad (5.11)$$

$$\mathfrak{h}_1^\varepsilon = \partial_2 (\widehat{\mathcal{H}}_2 \varphi^\varepsilon) + \partial_3 (\widehat{\mathcal{H}}_3 \varphi^\varepsilon) \quad \text{on } \omega_T. \quad (5.12)$$

REMARK The invertible part of the boundary matrix of a system allows to control the trace at the boundary of the so-called noncharacteristic component of the vector solution. Thus, with the system (5.8a) (whose boundary matrix is $-\mathcal{E}_{12}$, because of (2.28)) we have the control of $q^\varepsilon, u_1^\varepsilon$ at the boundary; therefore the components of \mathcal{U}^ε appearing in the boundary conditions (5.8c), (5.8d) are well defined.

The same holds true for (5.8b) where we can get the control of $\mathfrak{H}_2^\varepsilon, \mathfrak{H}_3^\varepsilon, \mathfrak{E}_2^\varepsilon, \mathfrak{E}_3^\varepsilon$. The control of $\mathfrak{E}_1^\varepsilon$ (which appears in (5.8d)) is not given from the system (5.8b), but from the constraint (5.10), as will be shown later on. We recall that $\mathfrak{H}_1^\varepsilon$ does not appear in the boundary condition (5.8d) because $\hat{\mathfrak{h}}_1 = \hat{\mathfrak{h}}_N = 0$.

Before studying problem (5.8) (or equivalently (5.1)), we should be sure that the number of boundary conditions is in agreement with the number of incoming characteristics for the hyperbolic systems (5.8). Since one of the four boundary conditions (5.8c)–(5.8f) is needed for determining the function $\varphi^\varepsilon(t, x')$, the total number of “incoming” characteristics should be three. Let us check that this is true.

PROPOSITION 5.1 If $0 < \varepsilon < 1$ system (5.8a) has one incoming characteristic for the boundary ω_T of the domain Q_T^+ . If $\varepsilon > 0$ is sufficiently small, system (5.8b) has two incoming characteristics for the boundary ω_T of the domain Q_T^- .

Proof. Consider first system (5.8a). In view of (2.28), the boundary matrix on ω_T is $-\mathcal{E}_{12}$ which has one negative (incoming in the domain Q_T^+) and one positive eigenvalue, while all other eigenvalues are zero.

Now consider system (5.8b). The boundary matrix B_1^ε has eigenvalues $\lambda_{1,2} = -\varepsilon^{-1}$, $\lambda_{3,4} = \varepsilon^{-1}$, $\lambda_{5,6} = 0$. Thus, system (5.8b) has indeed two incoming characteristics in the domain Q_T^- ($\lambda_{1,2} < 0$). \square

6. Basic a priori estimate for the hyperbolic regularized problem

Our goal now is to justify rigorously the formal limit $\varepsilon \rightarrow 0$ in (5.1)–(5.5), or alternatively in (5.8)–(5.12). To this end we will prove the existence of solutions to problem (5.8)–(5.12) and a uniform in ε a priori estimate. This work will be done in several steps.

6.1 The boundary value problem

Assuming that all coefficients and data appearing in (5.8) are extended for all times to the whole real line, let us consider the boundary value problem (recall the definition of Q^\pm, ω in (3.1))

$$\hat{\mathcal{Q}}_0 \partial_t \mathcal{U}^\varepsilon + \sum_{j=1}^3 (\hat{\mathcal{Q}}_j + \mathcal{E}_{1j+1}) \partial_j \mathcal{U}^\varepsilon + \hat{\mathcal{C}}' \mathcal{U}^\varepsilon = \mathfrak{F}, \quad \text{in } Q^+, \quad (6.1a)$$

$$B_0 \partial_t W^\varepsilon + \sum_{j=1}^3 B_j^\varepsilon \partial_j W^\varepsilon + B_4 W^\varepsilon = 0 \quad \text{in } Q^-, \quad (6.1b)$$

$$\partial_t \varphi^\varepsilon + \hat{v}_2 \partial_2 \varphi^\varepsilon + \hat{v}_3 \partial_3 \varphi^\varepsilon - \varphi^\varepsilon \partial_1 \hat{v}_N - u_1^\varepsilon = 0, \quad (6.1c)$$

$$q^\varepsilon + [\partial_1 \hat{q}] \varphi^\varepsilon - \hat{\mathfrak{h}} \cdot \mathfrak{H}^\varepsilon + \varepsilon \hat{\mathfrak{e}} \cdot \mathfrak{E}^\varepsilon = 0, \quad (6.1d)$$

$$\mathfrak{E}_2^\varepsilon - \varepsilon \partial_t (\hat{\mathcal{H}}_3 \varphi^\varepsilon) + \varepsilon \partial_2 (\hat{E}_1 \varphi^\varepsilon) = 0, \quad (6.1e)$$

$$\mathfrak{E}_3^\varepsilon + \varepsilon \partial_t (\widehat{\mathcal{H}}_2 \varphi^\varepsilon) + \varepsilon \partial_3 (\widehat{E}_1 \varphi^\varepsilon) = 0 \quad \text{on } \omega, \quad (6.1f)$$

$$(\mathcal{U}^\varepsilon, W^\varepsilon, \varphi^\varepsilon) = 0 \quad \text{for } t < 0. \quad (6.1g)$$

404 In this section we prove a uniform in ε a priori estimate of smooth solutions of (6.1).

THEOREM 6.1 Let the basic state (2.1) satisfies assumptions (2.2)–(2.7) and (4.1) for all times. There exist $\varepsilon_0 > 0$, $\gamma_0 \geq 1$ such that if $0 < \varepsilon < \varepsilon_0$ and $\gamma \geq \gamma_0$ then all sufficiently smooth solutions $(\mathcal{U}^\varepsilon, W^\varepsilon, \varphi^\varepsilon)$ of problem (6.1) obey the estimate

$$\begin{aligned} \gamma (\|\mathcal{U}_\gamma^\varepsilon\|_{H_{tan,\gamma}^1(Q^+)}^2 + \|W_\gamma^\varepsilon\|_{H_\gamma^1(Q^-)}^2 + \|(q_\gamma^\varepsilon, u_{1\gamma}^\varepsilon, h_{1\gamma}^\varepsilon)|_\omega\|_{H_\gamma^{1/2}(\omega)}^2 + \|W_\gamma^\varepsilon|_\omega\|_{H_\gamma^{1/2}(\omega)}^2) \\ + \gamma^2 \|\varphi_\gamma^\varepsilon\|_{H_\gamma^1(\omega)}^2 \leq \frac{C}{\gamma} \|\mathfrak{F}_\gamma\|_{H_{tan,\gamma}^1(Q^+)}^2, \end{aligned} \quad (6.2)$$

405 where we have set $\mathcal{U}_\gamma^\varepsilon = e^{-\gamma t} \mathcal{U}^\varepsilon$, $W_\gamma^\varepsilon = e^{-\gamma t} W^\varepsilon$, $\varphi_\gamma^\varepsilon = e^{-\gamma t} \varphi^\varepsilon$ and so on, and where $C =$
406 $C(K, \delta) > 0$ is a constant independent of the data \mathfrak{F} and the parameters ε, γ .

407 Passing to the limit $\varepsilon \rightarrow 0$ in this estimate will give the a priori estimate (4.2).

408 Since problem (6.1) looks similar to a corresponding one in relativistic MHD [24], for the
409 deduction of estimate (6.2) we use the same ideas as in [24]. On the one hand, we even have
410 an advantage, in comparison with the problem in [24], because the coefficients \widehat{E}_j in (6.1b),
411 (6.1d)–(6.1f) are still arbitrary functions whose choice will be crucial to make boundary conditions
412 dissipative. On the other hand, we should be more careful with lower-order terms than in [24],
413 because we must avoid the appearance of terms with ε^{-1} (otherwise, our estimate will not be
414 uniform in ε). Also for this reason we are using the variables $(\mathcal{U}^\varepsilon, W^\varepsilon)$ rather than $(U^\varepsilon, V^\varepsilon)$.

415 For the proof of (6.2) we will need a secondary symmetrization of the transformed Maxwell
416 equations in vacuum (5.1b), (5.3).

417 6.2 A secondary symmetrization

In order to show how to get the secondary symmetrization, for the sake of simplicity we consider first a planar unperturbed interface, i.e., the case $\widehat{\varphi} \equiv 0$. For this case (5.1b), (5.3) become

$$\partial_t V^\varepsilon + \sum_{j=1}^3 B_j^\varepsilon \partial_j V^\varepsilon = 0, \quad (6.3)$$

$$\operatorname{div} \mathcal{H}^\varepsilon = 0, \quad \operatorname{div} E^\varepsilon = 0. \quad (6.4)$$

418 We write for system (6.3) the following secondary symmetrization (for a similar secondary
419 symmetrization of the Maxwell equations in vacuum see [24]):

$$420 \quad \mathfrak{B}_0^\varepsilon \partial_t V^\varepsilon + \sum_{j=1}^3 \mathfrak{B}_j^\varepsilon B_j^\varepsilon \partial_j V^\varepsilon + R_1 \operatorname{div} \mathcal{H}^\varepsilon + R_2 \operatorname{div} E^\varepsilon = \mathfrak{B}_0^\varepsilon \partial_t V^\varepsilon + \sum_{j=1}^3 \mathfrak{B}_j^\varepsilon \partial_j V^\varepsilon = 0, \quad (6.5)$$

where

$$\mathfrak{B}_0^\varepsilon = \begin{pmatrix} 1 & 0 & 0 & 0 & \varepsilon\nu_3 & -\varepsilon\nu_2 \\ 0 & 1 & 0 & -\varepsilon\nu_3 & 0 & \varepsilon\nu_1 \\ 0 & 0 & 1 & \varepsilon\nu_2 & -\varepsilon\nu_1 & 0 \\ 0 & -\varepsilon\nu_3 & \varepsilon\nu_2 & 1 & 0 & 0 \\ \varepsilon\nu_3 & 0 & -\varepsilon\nu_1 & 0 & 1 & 0 \\ -\varepsilon\nu_2 & \varepsilon\nu_1 & 0 & 0 & 0 & 1 \end{pmatrix}, \quad (6.6)$$

$$\mathfrak{B}_1^\varepsilon = \begin{pmatrix} \nu_1 & \nu_2 & \nu_3 & 0 & 0 & 0 \\ \nu_2 & -\nu_1 & 0 & 0 & 0 & -\varepsilon^{-1} \\ \nu_3 & 0 & -\nu_1 & 0 & \varepsilon^{-1} & 0 \\ 0 & 0 & 0 & \nu_1 & \nu_2 & \nu_3 \\ 0 & 0 & \varepsilon^{-1} & \nu_2 & -\nu_1 & 0 \\ 0 & -\varepsilon^{-1} & 0 & \nu_3 & 0 & -\nu_1 \end{pmatrix}, \quad \mathfrak{B}_2^\varepsilon = \begin{pmatrix} -\nu_2 & \nu_1 & 0 & 0 & 0 & \varepsilon^{-1} \\ \nu_1 & \nu_2 & \nu_3 & 0 & 0 & 0 \\ 0 & \nu_3 & -\nu_2 & -\varepsilon^{-1} & 0 & 0 \\ 0 & 0 & -\varepsilon^{-1} & -\nu_2 & \nu_1 & 0 \\ 0 & 0 & 0 & \nu_1 & \nu_2 & \nu_3 \\ \varepsilon^{-1} & 0 & 0 & 0 & \nu_3 & -\nu_2 \end{pmatrix},$$

$$\mathfrak{B}_3^\varepsilon = \begin{pmatrix} -\nu_3 & 0 & \nu_1 & 0 & -\varepsilon^{-1} & 0 \\ 0 & -\nu_3 & \nu_2 & \varepsilon^{-1} & 0 & 0 \\ \nu_1 & \nu_2 & \nu_3 & 0 & 0 & 0 \\ 0 & \varepsilon^{-1} & 0 & -\nu_3 & 0 & \nu_1 \\ -\varepsilon^{-1} & 0 & 0 & 0 & -\nu_3 & \nu_2 \\ 0 & 0 & 0 & \nu_1 & \nu_2 & \nu_3 \end{pmatrix}, \quad R_1 = \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \\ 0 \\ 0 \\ 0 \end{pmatrix}, \quad R_2 = \begin{pmatrix} 0 \\ 0 \\ 0 \\ \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}.$$

421 The arbitrary functions $\nu_i(t, x)$ will be chosen in appropriate way later on. It may be useful to notice
422 that system (6.5) can also be written as

$$\begin{aligned} 423 \quad & (\partial_t \mathcal{H}^\varepsilon + \frac{1}{\varepsilon} \nabla \times E^\varepsilon) - \vec{v} \times (\varepsilon \partial_t E^\varepsilon - \nabla \times \mathcal{H}^\varepsilon) + \vec{v} \operatorname{div} \mathcal{H}^\varepsilon = 0, \\ & (\partial_t E^\varepsilon - \frac{1}{\varepsilon} \nabla \times \mathcal{H}^\varepsilon) + \vec{v} \times (\varepsilon \partial_t \mathcal{H}^\varepsilon + \nabla \times E^\varepsilon) + \vec{v} \operatorname{div} E^\varepsilon = 0, \end{aligned} \quad (6.7)$$

424 with the vector-function $\vec{v} = (\nu_1, \nu_2, \nu_3)$. The symmetric system (6.5) (or (6.7)) is hyperbolic if
425 $\mathfrak{B}_0^\varepsilon > 0$, i.e. for

$$426 \quad \varepsilon |\vec{v}| < 1. \quad (6.8)$$

427 The last inequality is satisfied for any given ν and small ε . We compute²

$$428 \quad \det(\mathfrak{B}_1^\varepsilon) = \nu_1^2 (|\vec{v}|^2 - 1/\varepsilon^2)^2.$$

429 Therefore the boundary is noncharacteristic for system (6.5) (or (6.7)) provided (6.8) and $\nu_1 \neq 0$
430 hold.

431 Consider now a nonplanar unperturbed interface, i.e., the general case when $\hat{\varphi}$ is not identically
432 zero. Similarly to (6.5), from (5.6), (5.3) we get the secondary symmetrization

$$433 \quad K \mathfrak{B}_0^\varepsilon K^{-1} \left(B_0 \partial_t W^\varepsilon + \sum_{j=1}^3 B_j^\varepsilon \partial_j W^\varepsilon + B_4 W^\varepsilon \right) + \frac{1}{\partial_1 \hat{\Phi}_1} K (R_1 \operatorname{div} \mathfrak{h}^\varepsilon + R_2 \operatorname{div} \mathfrak{e}^\varepsilon) = 0.$$

² The manual computation of the determinants is definitely too long. Here we used a free program for symbolic calculus, with the help of PS's son Martino.

We write this system as

$$M_0^\varepsilon \partial_t W^\varepsilon + \sum_{j=1}^3 M_j^\varepsilon \partial_j W^\varepsilon + M_4^\varepsilon W^\varepsilon = 0, \quad (6.9)$$

where

$$\begin{aligned} M_0^\varepsilon &= \frac{1}{\partial_1 \widehat{\Phi}_1} K \mathfrak{B}_0^\varepsilon K^\top > 0, \quad M_j^\varepsilon = \frac{1}{\partial_1 \widehat{\Phi}_1} K \mathfrak{B}_j^\varepsilon K^\top \quad (j = 2, 3), \\ M_1^\varepsilon &= \frac{1}{\partial_1 \widehat{\Phi}_1} K \widetilde{\mathfrak{B}}_1^\varepsilon K^\top, \quad \widetilde{\mathfrak{B}}_1^\varepsilon = \frac{1}{\partial_1 \widehat{\Phi}_1} \left(\mathfrak{B}_1^\varepsilon - \sum_{k=2}^3 \mathfrak{B}_k^\varepsilon \partial_k \widehat{\Psi} \right), \\ M_4^\varepsilon &= K \left(\mathfrak{B}_0^\varepsilon \partial_t + \widetilde{\mathfrak{B}}_1^\varepsilon \partial_1 + \mathfrak{B}_2^\varepsilon \partial_2 + \mathfrak{B}_3^\varepsilon \partial_3 + \mathfrak{B}_0^\varepsilon B_4 \right) \left(\frac{1}{\partial_1 \widehat{\Phi}_1} K^\top \right). \end{aligned} \quad (6.10)$$

System (6.9) is symmetric hyperbolic provided that (6.8) holds. We compute

$$\det(M_1^\varepsilon) = (1 + (\partial_2 \widehat{\Phi})^2 + (\partial_3 \widehat{\Phi})^2)^2 (v_1 - v_2 \partial_2 \widehat{\Phi} - v_3 \partial_3 \widehat{\Phi})^2 (|\vec{v}|^2 - 1/\varepsilon^2)^2, \quad (6.11)$$

and so the boundary is noncharacteristic for system (6.9) if and only if (6.8) holds and $v_1 \neq v_2 \partial_2 \widehat{\Phi} + v_3 \partial_3 \widehat{\Phi}$. System (6.9) originates from a linear combination of equations (5.1b) similar to (6.7), namely from

$$\begin{aligned} (\partial_t \mathfrak{h}^\varepsilon + \frac{1}{\varepsilon} \nabla \times \mathfrak{E}^\varepsilon) - \hat{\eta} (\vec{v} \times \hat{\eta}^{-1} (\varepsilon \partial_t \mathfrak{e}^\varepsilon - \nabla \times \mathfrak{H}^\varepsilon)) + \frac{\hat{\eta} \vec{v}}{\partial_1 \widehat{\Phi}_1} \operatorname{div} \mathfrak{h}^\varepsilon &= 0, \\ (\partial_t \mathfrak{e}^\varepsilon - \frac{1}{\varepsilon} \nabla \times \mathfrak{H}^\varepsilon) + \hat{\eta} (\vec{v} \times \hat{\eta}^{-1} (\varepsilon \partial_t \mathfrak{h}^\varepsilon + \nabla \times \mathfrak{E}^\varepsilon)) + \frac{\hat{\eta} \vec{v}}{\partial_1 \widehat{\Phi}_1} \operatorname{div} \mathfrak{e}^\varepsilon &= 0. \end{aligned} \quad (6.12)$$

We need to know which is the behavior of the above matrices in (6.10) w.r.t. ε as $\varepsilon \rightarrow 0$. In view of this, let us denote a generic matrix which is bounded w.r.t. ε by $O(1)$. Looking at (6.12) we immediately find

$$M_0^\varepsilon = O(1), \quad M_j^\varepsilon = B_j^\varepsilon + O(1) \quad (j = 1, 2, 3), \quad M_4^\varepsilon = O(1). \quad (6.13)$$

As the matrices M_0^ε and M_4^ε do not contain the multiplier ε^{-1} , their norms are bounded as $\varepsilon \rightarrow 0$. Recalling that the matrices B_j^ε are constant, we deduce as well that all the possible derivatives (with respect to t and x_j) of the matrices M_j^ε have bounded norms as $\varepsilon \rightarrow 0$.

6.3 Proof of Theorem 6.1

For the proof of our basic a priori estimate (6.2) we will apply the energy method to the symmetric hyperbolic systems (6.1a) and (6.9). In the sequel $\gamma_0 \geq 1$ denotes a generic constant sufficiently large which may increase from formula to formula, and C is a generic constant that may change from line to line.

First of all we provide some preparatory estimates. In particular, to estimate the weighted conormal derivative $Z_1 = \sigma \partial_1$ of \mathcal{U}^ε (recall the definition (3.3) of the γ -dependent norm of $H_{tan,\gamma}^1$) we do not need any boundary condition because the weight σ vanishes on ω . Applying to system (6.1a) the operator Z_1 and using standard arguments of the energy method,³ yields the

³ We multiply Z_1 (6.1a) by $e^{-\gamma t} Z_1 \mathcal{U}_\gamma^\varepsilon$ and integrate by parts over \mathcal{Q}^+ , then we use the Cauchy-Schwarz inequality.

460 inequality

$$461 \quad \gamma \|Z_1 \mathcal{U}_\gamma^\varepsilon\|_{L^2(Q^+)}^2 \leq \frac{C}{\gamma} \left\{ \|\mathcal{F}_\gamma\|_{H_{tan,\gamma}^1(Q^+)}^2 + \|\mathcal{U}_\gamma^\varepsilon\|_{H_{tan,\gamma}^1(Q^+)}^2 + \|\varepsilon_{12} \partial_1 \mathcal{U}_\gamma^\varepsilon\|_{L^2(Q^+)}^2 \right\}, \quad (6.14)$$

462 for $\gamma \geq \gamma_0$. On the other hand, directly from the equation (6.1a) we have

$$463 \quad \|\varepsilon_{12} \partial_1 \mathcal{U}_\gamma^\varepsilon\|_{L^2(Q^+)}^2 \leq C \left\{ \|\mathcal{F}_\gamma\|_{L^2(Q^+)}^2 + \|\mathcal{U}_\gamma^\varepsilon\|_{H_{tan,\gamma}^1(Q^+)}^2 \right\}, \quad (6.15)$$

464 where C is independent of ε, γ . Thus from (6.14), (6.15) we get

$$465 \quad \gamma \|Z_1 \mathcal{U}_\gamma^\varepsilon\|_{L^2(Q^+)}^2 \leq \frac{C}{\gamma} \left\{ \|\mathcal{F}_\gamma\|_{H_{tan,\gamma}^1(Q^+)}^2 + \|\mathcal{U}_\gamma^\varepsilon\|_{H_{tan,\gamma}^1(Q^+)}^2 \right\}, \quad \gamma \geq \gamma_0, \quad (6.16)$$

466 where C is independent of ε, γ . Furthermore, using the special structure of the boundary matrix in
 467 (6.1a) (see (2.28)) and the divergence constraint (5.9), we may estimate the normal derivative of the
 468 noncharacteristic part $\mathcal{U}_{n\gamma}^\varepsilon = e^{-\gamma t}(q^\varepsilon, u_1^\varepsilon, h_1^\varepsilon)$ of the “plasma” unknown $\mathcal{U}_\gamma^\varepsilon$:

$$469 \quad \|\partial_1 \mathcal{U}_{n\gamma}^\varepsilon\|_{L^2(Q^+)}^2 \leq C \left\{ \|\mathcal{F}_\gamma\|_{L^2(Q^+)}^2 + \|\mathcal{U}_\gamma^\varepsilon\|_{H_{tan,\gamma}^1(Q^+)}^2 \right\}, \quad (6.17)$$

470 where C is independent of ε, γ . In a similar way we wish to express the normal derivative of W^ε
 471 through its tangential derivatives. Here it is convenient to use system (6.1b) rather than (6.9). We
 472 multiply (6.1b) by ε and find from the obtained equation an explicit expression for the normal
 473 derivatives of $\mathfrak{H}_2^\varepsilon, \mathfrak{H}_3^\varepsilon, \mathfrak{E}_2^\varepsilon, \mathfrak{E}_3^\varepsilon$. An explicit expression for the normal derivatives of $\mathfrak{H}_1^\varepsilon, \mathfrak{E}_1^\varepsilon$ is found
 474 through the divergence constraints (5.10). Thus we can estimate the normal derivatives of all the
 475 components of W^ε through its tangential derivatives:

$$476 \quad \|\partial_1 W_\gamma^\varepsilon\|_{L^2(Q^-)}^2 \leq C \left\{ \gamma^2 \|W_\gamma^\varepsilon\|_{L^2(Q^-)}^2 + \|\partial_t W_\gamma^\varepsilon\|_{L^2(Q^-)}^2 + \sum_{k=2}^3 \|\partial_k W_\gamma^\varepsilon\|_{L^2(Q^-)}^2 \right\}, \quad (6.18)$$

477 where C does not depend on ε and γ , for all $\varepsilon \leq \varepsilon_0$.

478 As for the front function φ^ε we easily obtain from (6.1c) the L^2 estimate

$$479 \quad \gamma \|\varphi_\gamma^\varepsilon\|_{L^2(\omega)}^2 \leq \frac{C}{\gamma} \|u_{1\gamma}^\varepsilon\|_{L^2(\omega)}^2, \quad \gamma \geq \gamma_0, \quad (6.19)$$

480 where C is independent of γ . Furthermore, thanks to our basic assumption (4.1)⁴ we can resolve
 481 (5.11), (5.12) and (6.1c) for the space-time gradient $\nabla_{t,x'} \varphi_\gamma^\varepsilon = (\partial_t \varphi_\gamma^\varepsilon, \partial_2 \varphi_\gamma^\varepsilon, \partial_3 \varphi_\gamma^\varepsilon)$:

$$482 \quad \nabla_{t,x'} \varphi_\gamma^\varepsilon = \hat{a}_1 h_{1\gamma}^\varepsilon + \hat{a}_2 h_{1\gamma}^\varepsilon + \hat{a}_3 u_{1\gamma}^\varepsilon + \hat{a}_4 \varphi_\gamma^\varepsilon + \gamma \hat{a}_5 \varphi_\gamma^\varepsilon, \quad (6.20)$$

483 where the vector-functions $\hat{a}_\alpha = a_\alpha(\widehat{U}|_\omega, \widehat{\mathcal{H}}|_\omega)$ of coefficients can be easily written in explicit form.
 484 From (6.20) we get

$$485 \quad \|\nabla_{t,x'} \varphi_\gamma^\varepsilon\|_{L^2(\omega)} \leq C \left(\|\mathcal{U}_{n\gamma}^\varepsilon|_\omega\|_{L^2(\omega)} + \|W_\gamma^\varepsilon|_\omega\|_{L^2(\omega)} + \gamma \|\varphi_\gamma^\varepsilon\|_{L^2(\omega)} \right). \quad (6.21)$$

⁴ Under the conditions $\hat{H}_N = \hat{\mathcal{H}}_N = 0$ one has $|\hat{H} \times \hat{\mathcal{H}}|^2 = (\hat{H}_2 \hat{\mathcal{H}}_3 - \hat{H}_3 \hat{\mathcal{H}}_2)^2 (\nabla' \hat{\varphi})^2$ on ω , where we have set $(\nabla' \hat{\varphi}) := (1 + |\partial_2 \hat{\varphi}|^2 + |\partial_3 \hat{\varphi}|^2)^{1/2}$.

Now we prove a L^2 energy estimate for $(\mathcal{U}^\varepsilon, W^\varepsilon)$. We multiply (6.1a) by $e^{-\gamma t} \mathcal{U}_\gamma^\varepsilon$ and (6.9) by $e^{-\gamma t} W_\gamma^\varepsilon$, integrate by parts over Q^\pm , then we use the Cauchy-Schwarz inequality. We easily obtain

$$\begin{aligned} & \gamma \int_{Q^+} (\widehat{A}_0 \mathcal{U}_\gamma^\varepsilon, \mathcal{U}_\gamma^\varepsilon) dx dt + \gamma \int_{Q^-} (M_0^\varepsilon W_\gamma^\varepsilon, W_\gamma^\varepsilon) dx dt + \int_\omega \mathcal{Q}^\varepsilon dx' dt \\ & \leq C \left\{ \frac{1}{\gamma} \|\mathfrak{T}_\gamma\|_{L^2(Q^+)}^2 + \|\mathcal{U}_\gamma^\varepsilon\|_{L^2(Q^+)}^2 + \|W_\gamma^\varepsilon\|_{L^2(Q^-)}^2 \right\}, \quad (6.22) \end{aligned}$$

where we have denoted

$$\mathcal{Q}^\varepsilon = -\frac{1}{2}(\mathfrak{E}_{12} \mathcal{U}_\gamma^\varepsilon, \mathcal{U}_\gamma^\varepsilon)|_\omega + \frac{1}{2}(M_1^\varepsilon W_\gamma^\varepsilon, W_\gamma^\varepsilon)|_\omega.$$

Thanks to the properties of the matrices M_α^ε ($\alpha = \overline{0, 4}$) described in (6.13), the constant C in (6.22) is uniformly bounded in ε and γ . Let us calculate the quadratic form \mathcal{Q}^ε for the following choice of the functions v_j in the secondary symmetrization⁵:

$$v_1 = \hat{v}_2 \partial_2 \hat{\varphi} + \hat{v}_3 \partial_3 \hat{\varphi}, \quad v_k = \hat{v}_k, \quad k = 2, 3. \quad (6.23)$$

After long calculations we get (for simplicity we drop the index γ)

$$\mathcal{Q}^\varepsilon = -q^\varepsilon u_1^\varepsilon + \varepsilon^{-1}(\mathfrak{H}_3^\varepsilon \mathfrak{E}_2^\varepsilon - \mathfrak{H}_2^\varepsilon \mathfrak{E}_3^\varepsilon) + (\hat{v}_2 \mathfrak{H}_2^\varepsilon + \hat{v}_3 \mathfrak{H}_3^\varepsilon) \mathcal{H}_N^\varepsilon + (\hat{v}_2 \mathfrak{E}_2^\varepsilon + \hat{v}_3 \mathfrak{E}_3^\varepsilon) E_N^\varepsilon, \quad \text{on } \omega. \quad (6.24)$$

Now we insert the boundary conditions (5.12), (6.1c)–(6.1f) in the quadratic form \mathcal{Q}^ε , recalling also $\hat{\mathcal{H}}_N|_\omega = 0$ and noticing that

$$\hat{\mathbf{e}} \cdot \mathfrak{E}^\varepsilon = \widehat{E}_1 E_N^\varepsilon + \widehat{E}_{\tau_2} E_2^\varepsilon + \widehat{E}_{\tau_3} E_3^\varepsilon = \hat{\mathfrak{E}} \cdot \mathfrak{e}^\varepsilon.$$

Again after long calculations we get

$$\begin{aligned} \mathcal{Q}^\varepsilon = & (\widehat{E}_1 + \hat{v}_2 \widehat{\mathcal{H}}_3 - \hat{v}_3 \widehat{\mathcal{H}}_2) (\varepsilon E_N^\varepsilon \partial_t \varphi^\varepsilon + \mathfrak{H}_2^\varepsilon \partial_3 \varphi^\varepsilon - \mathfrak{H}_3^\varepsilon \partial_2 \varphi^\varepsilon) \\ & + (\varepsilon \widehat{E}_{\tau_2} E_2^\varepsilon + \varepsilon \widehat{E}_{\tau_3} E_3^\varepsilon) (\partial_t \varphi^\varepsilon + \hat{v}_2 \partial_2 \varphi^\varepsilon + \hat{v}_3 \partial_3 \varphi^\varepsilon) \\ & + \varphi^\varepsilon \{ -\gamma q^\varepsilon + [\partial_1 \hat{q}] u_1^\varepsilon - \partial_1 \hat{v}_N (q^\varepsilon + [\partial_1 \hat{q}] \varphi^\varepsilon) + (\gamma \widehat{\mathcal{H}}_3 + \partial_t \widehat{\mathcal{H}}_3 - \partial_2 \widehat{E}_1) (\mathfrak{H}_3^\varepsilon + \varepsilon \hat{v}_2 E_N^\varepsilon) \\ & + (\gamma \widehat{\mathcal{H}}_2 + \partial_t \widehat{\mathcal{H}}_2 + \partial_3 \widehat{E}_1) (\mathfrak{H}_2^\varepsilon - \varepsilon \hat{v}_3 E_N^\varepsilon) + (\partial_2 \widehat{\mathcal{H}}_2 + \partial_3 \widehat{\mathcal{H}}_3) (\hat{v}_2 \mathfrak{H}_2^\varepsilon + \hat{v}_3 \mathfrak{H}_3^\varepsilon) \} \quad \text{on } \omega. \quad (6.25) \end{aligned}$$

Thanks to the multiplicative factor ε in the boundary condition (6.1e), (6.1f), the critical term with the multiplier ε^{-1} in (6.24) has been dropped out. We make the following choice of the coefficients \widehat{E}_j in the boundary conditions (6.1d)–(6.1f):

$$\widehat{E} = \widehat{\mathcal{H}} \times \vec{v},$$

where \vec{v} is that of (6.23). For this choice

$$\widehat{E}_1 + \hat{v}_2 \widehat{\mathcal{H}}_3 - \hat{v}_3 \widehat{\mathcal{H}}_2 = 0, \quad \widehat{E}_{\tau_2} = 0, \quad \widehat{E}_{\tau_3} = 0, \quad (6.26)$$

⁵ Notice that the choice (6.23) makes the boundary characteristic, see (6.11).

and this leaves us with

$$\begin{aligned} \mathcal{Q}^\varepsilon = & \varphi^\varepsilon \{ -\gamma(q^\varepsilon + [\partial_1 \hat{q}]\varphi^\varepsilon) + [\partial_1 \hat{q}](u_1^\varepsilon + \varphi^\varepsilon \partial_1 \hat{v}_N) + \partial_1 \hat{v}_N q^\varepsilon \\ & + (\gamma \hat{\mathcal{H}}_3 + \partial_t \hat{\mathcal{H}}_3 - \partial_2 \hat{E}_1)(\mathfrak{H}_3^\varepsilon + \varepsilon \hat{v}_2 E_N^\varepsilon) + (\gamma \hat{\mathcal{H}}_2 + \partial_t \hat{\mathcal{H}}_2 + \partial_3 \hat{E}_1)(\mathfrak{H}_2^\varepsilon - \varepsilon \hat{v}_3 E_N^\varepsilon) \\ & + (\partial_2 \hat{\mathcal{H}}_2 + \partial_3 \hat{\mathcal{H}}_3)(\hat{v}_2 \mathfrak{H}_2^\varepsilon + \hat{v}_3 \mathfrak{H}_3^\varepsilon) \} \quad \text{on } \omega. \end{aligned}$$

499 Then we write in more convenient form the terms with coefficient γ substituting from (6.1d)

$$500 \quad -(q^\varepsilon + [\partial_1 \hat{q}]\varphi^\varepsilon) + \hat{\mathcal{H}}_2 \mathfrak{H}_2^\varepsilon + \hat{\mathcal{H}}_3 \mathfrak{H}_3^\varepsilon = \varepsilon \hat{\mathbf{e}} \cdot \mathfrak{E}^\varepsilon,$$

501 and we notice that

$$502 \quad \hat{\mathbf{e}} \cdot \mathfrak{E}^\varepsilon + (\hat{v}_2 \hat{\mathcal{H}}_3 - \hat{v}_3 \hat{\mathcal{H}}_2) E_N^\varepsilon = \hat{\mathfrak{E}} \cdot \mathfrak{e}^\varepsilon + (\hat{v}_2 \hat{\mathcal{H}}_3 - \hat{v}_3 \hat{\mathcal{H}}_2) \mathfrak{e}_1^\varepsilon = 0, \quad \text{on } \omega,$$

again by (6.26). Thus we get

$$\begin{aligned} \mathcal{Q}^\varepsilon = & \varphi^\varepsilon \{ [\partial_1 \hat{q}](u_1^\varepsilon + \varphi^\varepsilon \partial_1 \hat{v}_N) + \partial_1 \hat{v}_N q^\varepsilon \\ & + (\partial_t \hat{\mathcal{H}}_3 - \partial_2 \hat{E}_1)(\mathfrak{H}_3^\varepsilon + \varepsilon \hat{v}_2 E_N^\varepsilon) + (\partial_t \hat{\mathcal{H}}_2 + \partial_3 \hat{E}_1)(\mathfrak{H}_2^\varepsilon - \varepsilon \hat{v}_3 E_N^\varepsilon) \\ & + (\partial_2 \hat{\mathcal{H}}_2 + \partial_3 \hat{\mathcal{H}}_3)(\hat{v}_2 \mathfrak{H}_2^\varepsilon + \hat{v}_3 \mathfrak{H}_3^\varepsilon) \} \quad \text{on } \omega. \quad (6.27) \end{aligned}$$

From (6.22), (6.27) we obtain (we restore the index γ)

$$\begin{aligned} \gamma \left(\|\mathfrak{U}_\gamma^\varepsilon\|_{L^2(Q^+)}^2 + \|W_\gamma^\varepsilon\|_{L^2(Q^-)}^2 \right) & \leq \frac{C}{\gamma} \left\{ \|\mathfrak{F}_\gamma^\varepsilon\|_{L^2(Q^+)}^2 + \|\mathfrak{U}_{n\gamma}^\varepsilon|_\omega\|_{L^2(\omega)}^2 + \|W_\gamma^\varepsilon|_\omega\|_{L^2(\omega)}^2 \right\} \\ & + C \left(\|\mathfrak{U}_\gamma^\varepsilon\|_{L^2(Q^+)}^2 + \|W_\gamma^\varepsilon\|_{L^2(Q^-)}^2 \right) + \gamma \|\varphi^\varepsilon\|_{L^2(\omega)}^2, \quad (6.28) \end{aligned}$$

where C is independent of ε, γ . Thus if γ_0 is large enough we obtain from (6.19), (6.28) the inequality

$$\begin{aligned} \gamma \left(\|\mathfrak{U}_\gamma^\varepsilon\|_{L^2(Q^+)}^2 + \|W_\gamma^\varepsilon\|_{L^2(Q^-)}^2 \right) & \leq \frac{C}{\gamma} \left\{ \|\mathfrak{F}_\gamma^\varepsilon\|_{L^2(Q^+)}^2 + \|\mathfrak{U}_{n\gamma}^\varepsilon|_\omega\|_{L^2(\omega)}^2 + \|W_\gamma^\varepsilon|_\omega\|_{L^2(\omega)}^2 \right\}, \quad 0 < \varepsilon < \varepsilon_0, \quad \gamma \geq \gamma_0, \quad (6.29) \end{aligned}$$

503 where C is independent of ε, γ .

Now we derive the a priori estimate of tangential derivatives. Differentiating systems (6.1a) and (6.9) with respect to $x_0 = t, x_2$ or x_3 , using standard arguments of the energy method, and applying (6.17), (6.18), gives the energy inequality

$$\begin{aligned} \gamma \int_{Q^+} (\hat{A}_0 Z_\ell U_\gamma^\varepsilon, Z_\ell U_\gamma^\varepsilon) dx dt + \gamma \int_{Q^-} (M_0^\varepsilon Z_\ell W_\gamma^\varepsilon, Z_\ell W_\gamma^\varepsilon) dx dt + \int_\omega \mathcal{Q}_\ell^\varepsilon dx' dt \\ \leq \frac{C}{\gamma} \left\{ \|\mathfrak{F}_\gamma^\varepsilon\|_{H_{tan,\gamma}^1(Q^+)}^2 + \|\mathfrak{U}_\gamma^\varepsilon\|_{H_{tan,\gamma}^1(Q^+)}^2 + \|W_\gamma^\varepsilon\|_{H_\gamma^1(Q^-)}^2 \right\}, \quad (6.30) \end{aligned}$$

where $\ell = 0, 2, 3$, and where we have denoted

$$\mathcal{Q}_\ell^\varepsilon = -\frac{1}{2}(\mathfrak{E}_{12} Z_\ell U_\gamma^\varepsilon, Z_\ell U_\gamma^\varepsilon)|_\omega + \frac{1}{2}(M_1^\varepsilon Z_\ell W_\gamma^\varepsilon, Z_\ell W_\gamma^\varepsilon)|_\omega.$$

Thanks to the properties of the matrices M_α^ε ($\alpha = \overline{0, 4}$) described in (6.13), the constant C in (6.30) is uniformly bounded in ε and γ . We repeat for $\mathcal{Q}_\ell^\varepsilon$ the calculations leading to (6.27) for \mathcal{Q}^ε . Clearly, for the same choices as in (6.23) and (6.26) we obtain (for simplicity we drop again the index γ)

$$\begin{aligned} \mathcal{Q}_\ell^\varepsilon = & Z_\ell \varphi^\varepsilon \{ [\partial_1 \hat{q}] (Z_\ell u_1^\varepsilon + Z_\ell \varphi^\varepsilon \partial_1 \hat{v}_N) + \partial_1 \hat{v}_N Z_\ell \varphi^\varepsilon \\ & + (\partial_t \hat{\mathcal{H}}_3 - \partial_2 \hat{E}_1)(Z_\ell \mathfrak{H}_3^\varepsilon + \varepsilon \hat{v}_2 Z_\ell E_N^\varepsilon) + (\partial_t \hat{\mathcal{H}}_2 + \partial_3 \hat{E}_1)(Z_\ell \mathfrak{H}_2^\varepsilon - \varepsilon \hat{v}_3 Z_\ell E_N^\varepsilon) \\ & + (\partial_2 \hat{\mathcal{H}}_2 + \partial_3 \hat{\mathcal{H}}_3)(\hat{v}_2 Z_\ell \mathfrak{H}_2^\varepsilon + \hat{v}_3 Z_\ell \mathfrak{H}_3^\varepsilon) \} + \text{l.o.t.}, \quad \text{on } \omega, \end{aligned} \quad (6.31)$$

where l.o.t. is the sum of lower-order terms. Using (6.20) we reduce the above terms to those like

$$\hat{c} h_1^\varepsilon Z_\ell u_1^\varepsilon, \quad \hat{c} h_1^\varepsilon Z_\ell \varphi^\varepsilon, \quad \hat{c} h_1^\varepsilon Z_\ell \mathfrak{H}_j^\varepsilon, \quad \hat{c} h_1^\varepsilon Z_\ell \mathfrak{E}_j^\varepsilon, \quad \dots \quad \text{on } \omega,$$

terms as above with $\mathfrak{h}_1^\varepsilon, u_1^\varepsilon$ instead of h_1^ε , or even “better” terms like

$$\gamma \hat{c} \varphi^\varepsilon Z_\ell u_1^\varepsilon, \quad \gamma \hat{c} \varphi^\varepsilon Z_\ell \varphi^\varepsilon.$$

Here and below \hat{c} is the common notation for a generic coefficient depending on the basic state (2.1). By integration by parts such “better” terms can be reduced to the above ones and terms of lower order.

The terms like $\hat{c} h_1^\varepsilon Z_\ell u_{1|x_1=0}^\varepsilon$ are estimated by passing to the volume integral and integrating by parts:

$$\begin{aligned} & \int_\omega \hat{c} h_1^\varepsilon Z_\ell u_{1|x_1=0}^\varepsilon dx' dt \\ &= - \int_{Q^+} \partial_1 (\tilde{c} h_1^\varepsilon Z_\ell u_1^\varepsilon) dx dt \\ &= \int_{Q^+} \left\{ (Z_\ell \tilde{c}) h_1^\varepsilon (\partial_1 u_1^\varepsilon) + \tilde{c} (Z_\ell h_1^\varepsilon) \partial_1 u_1^\varepsilon - (\partial_1 \tilde{c}) h_1^\varepsilon Z_\ell u_1^\varepsilon - \tilde{c} (\partial_1 h_1^\varepsilon) Z_\ell u_1^\varepsilon \right\} dx dt, \end{aligned}$$

where $\tilde{c}|_{x_1=0} = \hat{c}$. Estimating the right-hand side by the Hölder’s inequality and (6.17) gives

$$\left| \int_\omega \hat{c} h_1^\varepsilon Z_\ell u_{1|x_1=0}^\varepsilon dx' dt \right| \leq C \left\{ \|\mathfrak{T}_\gamma\|_{L^2(Q^+)}^2 + \|\mathcal{U}_\gamma^\varepsilon\|_{H_{tan,\gamma}^1(Q^+)}^2 \right\}. \quad (6.32)$$

In the same way we estimate the other similar terms $\hat{c} h_1^\varepsilon Z_\ell \mathfrak{H}_j^\varepsilon, \hat{c} h_1^\varepsilon Z_\ell \mathfrak{E}_j^\varepsilon$, etc. Notice that we only need to estimate normal derivatives either of components of $\mathcal{U}_{n\gamma}^\varepsilon$ or W_γ^ε . For terms like $\hat{c} \mathfrak{h}_1^\varepsilon Z_\ell u_1^\varepsilon, \hat{c} \mathfrak{h}_1^\varepsilon Z_\ell \mathfrak{E}_j^\varepsilon$, etc. we use (6.18) instead of (6.17).

We treat the terms like $\hat{c} h_{1|x_1=0}^\varepsilon Z_\ell \varphi^\varepsilon$ by substituting (6.20) again:

$$\begin{aligned} \left| \int_\omega \hat{c} h_1^\varepsilon Z_\ell \varphi^\varepsilon dx' dt \right| &= \left| \int_\omega \hat{c} h_1^\varepsilon (\hat{a}_1 h_1^\varepsilon + \hat{a}_2 \mathfrak{h}_1^\varepsilon + \hat{a}_3 u_1^\varepsilon + \hat{a}_4 \varphi^\varepsilon + \gamma \hat{a}_5 \varphi^\varepsilon) dx dt \right| \\ &\leq C \left(\|\mathcal{U}_n^\varepsilon\|_{L^2(\omega)}^2 + \|W^\varepsilon\|_{L^2(\omega)}^2 + \gamma^2 \|\varphi^\varepsilon\|_{L^2(\omega)}^2 \right). \end{aligned} \quad (6.33)$$

Combining (6.30), (6.32), (6.33) and similar inequalities for the other terms of (6.31) yields (we

restore the index γ)

$$\begin{aligned} & \gamma \left(\|Z_\ell \mathcal{U}_\gamma^\varepsilon\|_{L^2(Q^+)}^2 + \|Z_\ell W_\gamma^\varepsilon\|_{L^2(Q^-)}^2 \right) \\ & \leq C \left\{ \frac{1}{\gamma} \|\mathfrak{F}_\gamma\|_{H_{tan,\gamma}^1(Q^+)}^2 + \|\mathcal{U}_\gamma^\varepsilon\|_{H_{tan,\gamma}^1(Q^+)}^2 + \|W_\gamma^\varepsilon\|_{H_\gamma^1(Q^-)}^2 \right. \\ & \quad \left. + \gamma \left(\|\mathcal{U}_{n\gamma}^\varepsilon\|_{L^2(\omega)}^2 + \|W_\gamma^\varepsilon\|_{L^2(\omega)}^2 \right) \right\}, \quad 0 < \varepsilon < \varepsilon_0, \gamma \geq \gamma_0, \end{aligned} \quad (6.34)$$

where C is independent of ε, γ . Then from (6.16), (6.18), (6.29), (6.34) we obtain

$$\begin{aligned} & \gamma \left(\|\mathcal{U}_\gamma^\varepsilon\|_{H_{tan,\gamma}^1(Q^+)}^2 + \|W_\gamma^\varepsilon\|_{H_\gamma^1(Q^-)}^2 \right) \\ & \leq C \left\{ \frac{1}{\gamma} \|\mathfrak{F}_\gamma\|_{H_{tan,\gamma}^1(Q^+)}^2 + \|\mathcal{U}_\gamma^\varepsilon\|_{H_{tan,\gamma}^1(Q^+)}^2 + \|W_\gamma^\varepsilon\|_{H_\gamma^1(Q^-)}^2 \right. \\ & \quad \left. + \gamma \left(\|\mathcal{U}_{n\gamma}^\varepsilon\|_{L^2(\omega)}^2 + \|W_\gamma^\varepsilon\|_{L^2(\omega)}^2 \right) \right\}, \quad 0 < \varepsilon < \varepsilon_0, \gamma \geq \gamma_0, \end{aligned} \quad (6.35)$$

516 where C is independent of ε, γ . We need the following estimates for the trace of $\mathcal{U}_n^\varepsilon, W^\varepsilon$.

LEMMA 6.2 The functions $\mathcal{U}_n^\varepsilon, W^\varepsilon$ satisfy

$$\gamma \|\mathcal{U}_{n\gamma}^\varepsilon\|_{L^2(\omega)}^2 + \|\mathcal{U}_{n\gamma}^\varepsilon\|_{H_\gamma^{1/2}(\omega)}^2 \leq C \left(\|\mathfrak{F}_\gamma\|_{L^2(Q^+)}^2 + \|\mathcal{U}_\gamma^\varepsilon\|_{H_{tan,\gamma}^1(Q^+)}^2 \right), \quad (6.36)$$

$$\gamma \|W_\gamma^\varepsilon\|_{L^2(\omega)}^2 + \|W_\gamma^\varepsilon\|_{H_\gamma^{1/2}(\omega)}^2 \leq C \|W_\gamma^\varepsilon\|_{H_\gamma^1(Q^-)}^2. \quad (6.37)$$

517 The proof of Lemma 6.2 is given in Section 11 at the end of this article. Substituting (6.36),
518 (6.37) in (6.35) and taking γ_0 large enough yields

$$519 \quad \gamma \left(\|\mathcal{U}_\gamma^\varepsilon\|_{H_{tan,\gamma}^1(Q^+)}^2 + \|W_\gamma^\varepsilon\|_{H_\gamma^1(Q^-)}^2 \right) \leq \frac{C}{\gamma} \|\mathfrak{F}_\gamma\|_{H_{tan,\gamma}^1(Q^+)}^2, \quad 0 < \varepsilon < \varepsilon_0, \gamma \geq \gamma_0, \quad (6.38)$$

520 where C is independent of ε, γ . Finally, from (6.21), (6.36) and (6.38) we get

$$521 \quad \gamma \left(\|\mathcal{U}_{n\gamma}^\varepsilon\|_{H_\gamma^{1/2}(\omega)}^2 + \|W_\gamma^\varepsilon\|_{H_\gamma^{1/2}(\omega)}^2 \right) + \gamma^2 \|\varphi^\varepsilon\|_{H_\gamma^1(\omega)}^2 \leq \frac{C}{\gamma} \|\mathfrak{F}_\gamma\|_{H_{tan,\gamma}^1(Q^+)}^2. \quad (6.39)$$

522 Adding (6.38), (6.39) gives (6.2). The proof of Theorem 6.1 is complete.

523 7. Well-posedness of the hyperbolic regularized problem

524 In this section we prove the existence of the solution of (6.1). Its restriction to the time interval
525 $(-\infty, T]$ will provide the solution of problem (5.8). From now on, in the proof of the existence of
526 the solution, ε is fixed and so we omit it and we simply write \mathcal{U} instead of \mathcal{U}^ε , W instead of W^ε , φ
527 instead of φ^ε .

528 In view of the result of Lemma 9.1 (see Section 9) we can consider system (6.9) instead of (6.1b).
529 First of all, we write the boundary conditions in different form, by eliminating the derivatives of φ .

530 We substitute (6.1c) in the boundary conditions for $\mathfrak{E}_2, \mathfrak{E}_3$ and take account of the constraint (5.12)
 531 and the choices (6.23), (6.26). We get

$$\begin{aligned} q - \hat{\mathfrak{h}}_2 \mathfrak{H}_2 - \hat{\mathfrak{h}}_3 \mathfrak{H}_3 + \varepsilon \hat{E}_1 E_N + [\partial_1 \hat{q}] \varphi &= 0, \\ \mathfrak{E}_2 - \varepsilon \hat{\mathcal{H}}_3 u_1 + \varepsilon \hat{v}_3 \mathcal{H}_N + \varepsilon a_1 \varphi &= 0, \\ \mathfrak{E}_3 + \varepsilon \hat{\mathcal{H}}_2 u_1 - \varepsilon \hat{v}_2 \mathcal{H}_N + \varepsilon a_2 \varphi &= 0, \quad \text{on } \omega, \end{aligned} \quad (7.1)$$

533 where the precise form of the coefficients a_1, a_2 is not important. For later use we observe that
 534 (5.12), (6.1c)–(6.1f) is equivalent to (5.12), (6.1c), (7.1). Notice that the last two equations in (7.1)
 535 yield

$$\varepsilon \hat{E}_1 u_1 + \hat{v}_2 \mathfrak{E}_2 + \hat{v}_3 \mathfrak{E}_3 + \varepsilon a_3 \varphi = 0, \quad (7.2)$$

537 where $a_3 = a_1 \hat{v}_2 + a_2 \hat{v}_3$.

538 Let us write the system (6.1a), (6.9), (7.1) in compact form as

$$\begin{cases} \mathfrak{L} \begin{pmatrix} \mathcal{U} \\ W \end{pmatrix} = \begin{pmatrix} \mathfrak{F} \\ 0 \end{pmatrix} & \text{on } Q^+ \times Q^-, \\ M \begin{pmatrix} \mathcal{U} \\ W \end{pmatrix} + b \varphi = 0, & \text{in } \omega, \\ (\mathcal{U}, W, \varphi) = 0 & \text{for } t < 0, \end{cases} \quad (7.3)$$

540 where the matrix M and the vector b are implicitly defined by (7.1).

541 Let us multiply (7.3) by $e^{-\gamma t}$ with $\gamma \geq 1$; according to the rule $e^{-\gamma t} \partial_t u = (\gamma + \partial_t) e^{-\gamma t} u$, (7.3)
 542 becomes equivalent to

$$\begin{cases} \mathfrak{L}_\gamma \begin{pmatrix} \mathcal{U}_\gamma \\ W_\gamma \end{pmatrix} = \begin{pmatrix} \mathfrak{F}_\gamma \\ 0 \end{pmatrix} & \text{on } Q^+ \times Q^-, \\ M \begin{pmatrix} \mathcal{U}_\gamma \\ W_\gamma \end{pmatrix} + b \varphi_\gamma = 0 & \text{in } \omega, \\ (\mathcal{U}_\gamma, W_\gamma, \varphi_\gamma) = 0 & \text{for } t < 0. \end{cases} \quad (7.4)$$

544 where

$$\mathfrak{L}_\gamma := \gamma \begin{pmatrix} \hat{\mathfrak{G}}_0 & 0 \\ 0 & M_0^\varepsilon \end{pmatrix} + \mathfrak{L},$$

546 $\mathcal{U}_\gamma = e^{-\gamma t} \mathcal{U}, W_\gamma = e^{-\gamma t} W, \varphi_\gamma = e^{-\gamma t} \varphi$, etc.

547 First we solve (7.4) under the assumption that φ_γ is given.

LEMMA 7.1 There exists $\gamma_0 > 0$ such that for all $\gamma \geq \gamma_0$ and for all given $\mathfrak{F} \in e^{\gamma t} H_{tan, \gamma}^1(Q^+)$
 and $\varphi \in e^{\gamma t} H_\gamma^{3/2}(\omega)$ vanishing in the past, the problem (7.4) has a unique solution $(\mathcal{U}, W) \in$

$e^{\gamma t} H_{tan,\gamma}^1(Q^+) \times e^{\gamma t} H_\gamma^1(Q^-)$ with $(q, u_1 h_1, W)|_\omega \in e^{\gamma t} H_\gamma^{1/2}(\omega)$, such that

$$\begin{aligned} & \|e^{-\gamma t} (\mathcal{U}, W)\|_{H_{tan,\gamma}^1(Q^+) \times H_\gamma^1(Q^-)} + \|e^{-\gamma t} (q, u_1, h_1, W)|_\omega\|_{H_\gamma^{1/2}(\omega)} \\ & \leq \frac{C}{\gamma} \left(\|e^{-\gamma t} \mathfrak{F}\|_{H_{tan,\gamma}^1(Q^+)} + \|e^{-\gamma t} \varphi\|_{H_\gamma^{3/2}(\omega)} \right). \end{aligned} \quad (7.5)$$

548 *Proof.* We insert the new boundary conditions (7.1), (7.2) in the quadratic form \mathcal{Q}^ε (see (6.24)) and
549 we get

$$550 \quad \mathcal{Q}^\varepsilon := -\frac{1}{2}(\hat{\mathcal{Q}}_1 + \mathcal{E}_{12}) \mathcal{U} \cdot \mathcal{U} + \frac{1}{2} M_1^\varepsilon W \cdot W = ([\partial \hat{q}] u_1 + a_2 \mathfrak{H}_2 - a_1 \mathfrak{H}_3 - \varepsilon a_3 E_N) \varphi \quad \text{on } \omega. \quad (7.6)$$

551 If we consider the boundary conditions (7.1), (7.2) in homogeneous form, namely if we set $\varphi = 0$,
552 then from (7.6)

$$553 \quad \mathcal{Q}^\varepsilon = 0 \quad \text{on } \omega.$$

We deduce that the boundary conditions (7.1) are nonnegative for \mathfrak{L}_γ . As the number of boundary conditions in (7.1) is in agreement with the number of incoming characteristics for the operator \mathfrak{L}_γ (see Proposition 5.1) we infer that the boundary conditions (7.1) are maximally nonnegative (but not strictly dissipative). Then we reduce the problem to one with homogeneous boundary conditions by subtracting from $(\mathcal{U}_\gamma, W_\gamma)$ a function $(\mathcal{U}'_\gamma, W'_\gamma) \in H_\gamma^2(Q^+) \times H_\gamma^2(Q^-)$ such that

$$M \begin{pmatrix} \mathcal{U}'_\gamma \\ W'_\gamma \end{pmatrix} + b \varphi = 0 \quad \text{on } \omega.$$

554 Finally, as the boundary is characteristic of constant multiplicity [18], we may apply the result
555 of [19, 20] and we get the solution with the prescribed regularity. \square

556 The well-posedness of (6.1) in $H_{tan}^1 \times H^1$ is given by the following theorem.

557 **THEOREM 7.2** There exists $\gamma_0 > 0$ such that for all $\gamma \geq \gamma_0$ and $\mathfrak{F} \in e^{\gamma t} H_{tan,\gamma}^1(Q^+)$ vanishing
558 in the past, the problem (6.1) has a unique solution $(\mathcal{U}, W) \in e^{\gamma t} H_{tan,\gamma}^1(Q^+) \times e^{\gamma t} H_\gamma^1(Q^-)$ with
559 $(q, u_1 h_1, W)|_\omega \in e^{\gamma t} H_\gamma^{1/2}(\omega)$, $\varphi \in e^{\gamma t} H_\gamma^{3/2}(\omega)$.

560 *Proof.* We prove the existence of the solution to (6.1) by a fixed point argument. Let $\bar{\varphi} \in$
561 $e^{\gamma t} H_\gamma^{3/2}(\omega_T)$ vanishing in the past. By Lemma 7.1, for γ sufficiently large there exists a unique
562 solution $(\mathcal{U}, W) \in e^{\gamma t} H_{tan,\gamma}^1(Q^+) \times e^{\gamma t} H_\gamma^1(Q^-)$, with $(q, u_1, h_1, W)|_\omega \in e^{\gamma t} H_\gamma^{1/2}(\omega)$ of

$$563 \quad \begin{cases} \mathfrak{L}_\gamma \begin{pmatrix} \mathcal{U}_\gamma \\ W_\gamma \end{pmatrix} = \begin{pmatrix} \mathfrak{F}_\gamma \\ 0 \end{pmatrix} & \text{on } Q^+ \times Q^-, \\ M \begin{pmatrix} \mathcal{U}_\gamma \\ W_\gamma \end{pmatrix} = -b \bar{\varphi}_\gamma & \text{on } \omega, \\ (\mathcal{U}_\gamma, W_\gamma) = 0 & \text{for } t < 0, \end{cases} \quad (7.7)$$

564 enjoying the a priori estimate (7.5) with $\bar{\varphi}$ instead of φ . Now consider the equation

$$565 \quad \gamma \varphi_\gamma + \partial_t \varphi_\gamma + \hat{v}_2 \partial_2 \varphi_\gamma + \hat{v}_3 \partial_3 \varphi_\gamma - \varphi_\gamma \partial_1 \hat{v}_N = u_{1\gamma}, \quad \text{on } \omega, \quad (7.8)$$

where $u_{1\gamma} \in H_\gamma^{1/2}(\omega)$ is the trace of the component of \mathcal{U}_γ given in the previous step, vanishing for $t < 0$. For γ sufficiently large there exists a unique solution $\varphi_\gamma \in H_\gamma^{1/2}(\omega)$, vanishing in the past, such that

$$\|\varphi_\gamma\|_{H_\gamma^{1/2}(\omega)} \leq \frac{C}{\gamma} \|u_{1\gamma}\|_{H_\gamma^{1/2}(\omega)}. \quad (7.9)$$

From the plasma equation in (7.7) and from (7.8) we deduce the boundary constraint

$$h_{1\gamma} = \widehat{H}_2 \partial_2 \varphi_\gamma + \widehat{H}_3 \partial_3 \varphi_\gamma - \varphi_\gamma \partial_1 \widehat{H}_N \quad \text{on } \omega. \quad (7.10)$$

Since in the right-hand side of (7.7) we have $\overline{\varphi}$ instead of φ we are not able to deduce the similar boundary constraint for the vacuum magnetic field. Instead, we obtain

$$\mathfrak{h}_{1\gamma} - \partial_2(\widehat{\mathcal{H}}_2 \varphi_\gamma) - \partial_3(\widehat{\mathcal{H}}_3 \varphi_\gamma) = G_\gamma \quad \text{on } \omega, \quad (7.11)$$

where G_γ solves

$$\frac{\tilde{d}}{dt} G_\gamma + a_2 \partial_2 (\varphi_\gamma - \overline{\varphi}_\gamma) - a_1 \partial_3 (\varphi_\gamma - \overline{\varphi}_\gamma) + (\partial_2 a_2 - \partial_3 a_1) (\varphi_\gamma - \overline{\varphi}_\gamma) = 0 \quad \text{on } \omega, \quad (7.12)$$

for $\tilde{d}/dt = \gamma + \partial_t + \partial_2(\widehat{v}_2 \cdot) + \partial_3(\widehat{v}_3 \cdot)$ and where the coefficients a_1, a_2 are the same of (7.1). (7.12) is derived from the first equation of the vacuum part of (7.7), (7.8) and the boundary conditions for $\mathfrak{E}_2, \mathfrak{E}_3$ in (7.7).

Let us consider the linear system for $\nabla_{t,x'} \varphi_\gamma$ provided by equations (7.8), (7.10) and (7.11). By the stability condition (4.1) we can express $\nabla_{t,x'} \varphi_\gamma$ through $(h_{1\gamma}, \mathfrak{h}_{1\gamma}, u_{1\gamma})|_\omega, \varphi_\gamma, G_\gamma$, that is

$$\nabla_{t,x'} \varphi_\gamma = a'_1 h_{1\gamma} + a'_2 \mathfrak{h}_{1\gamma} + a'_3 u_{1\gamma} + a'_4 \varphi_\gamma + a'_5 G_\gamma, \quad (7.13)$$

where the precise form of the coefficients a'_i has no interest. Then, substituting into (7.12) yields

$$\frac{\tilde{d}}{dt} G_\gamma + b_0 G_\gamma = b_1 h_{1\gamma} + b_2 \mathfrak{h}_{1\gamma} + b_3 \varphi_\gamma + a_2 \partial_2 \overline{\varphi}_\gamma - a_1 \partial_3 \overline{\varphi}_\gamma + (\partial_2 a_2 - \partial_3 a_1) \overline{\varphi}_\gamma \quad \text{on } \omega, \quad (7.14)$$

with suitable coefficients b_i .

From (7.14), for γ sufficiently large, we get the estimate

$$\begin{aligned} \|G_\gamma\|_{H_\gamma^{1/2}(\omega)} &\leq \frac{C}{\gamma} \left(\|(h_{1\gamma}, \mathfrak{h}_{1\gamma})\|_{H_\gamma^{1/2}(\omega)} + \|\varphi_\gamma\|_{H_\gamma^{1/2}(\omega)} + \|\overline{\varphi}_\gamma\|_{H_\gamma^{3/2}(\omega)} \right) \\ &\leq \frac{C}{\gamma} \left(\|\mathfrak{F}_\gamma\|_{H_{tan,\gamma}^1(Q^+)} + \|\overline{\varphi}_\gamma\|_{H_\gamma^{3/2}(\omega)} \right), \end{aligned} \quad (7.15)$$

where we have applied (7.5) (with $\overline{\varphi}$ in place of φ) and (7.9). Thus, from (7.13) again, we obtain the estimate

$$\begin{aligned} \|\nabla_{t,x'} \varphi_\gamma\|_{H_\gamma^{1/2}(\omega)} &\leq C \left(\|(u_{1\gamma}, h_{1\gamma}, \mathfrak{h}_{1\gamma})\|_{H_\gamma^{1/2}(\omega)} + \|\varphi_\gamma\|_{H_\gamma^{1/2}(\omega)} + \|G_\gamma\|_{H_\gamma^{1/2}(\omega)} \right) \\ &\leq \frac{C}{\gamma} \left(\|\mathfrak{F}_\gamma\|_{H_{tan,\gamma}^1(Q^+)} + \|\overline{\varphi}_\gamma\|_{H_\gamma^{3/2}(\omega)} \right). \end{aligned} \quad (7.16)$$

Combining (7.5) (with $\bar{\varphi}$ in place of φ), (7.9) and (7.16) gives

$$\|\varphi_\gamma\|_{H_\gamma^{3/2}(\omega)} \leq \frac{C}{\gamma} \left(\|\mathfrak{F}_\gamma\|_{H_{tan,\gamma}^1(Q^+)} + \|\bar{\varphi}_\gamma\|_{H_\gamma^{3/2}(\omega)} \right). \quad (7.17)$$

This defines a map $\bar{\varphi} \rightarrow \varphi$ in $e^{\gamma t} H_\gamma^{3/2}(\omega_T)$. Let $\bar{\varphi}^1, \bar{\varphi}^2 \in e^{\gamma t} H_\gamma^{3/2}(\omega_T)$, and $(\mathcal{U}^1, W^1), (\mathcal{U}^2, W^2)$, φ^1, φ^2 be the corresponding solutions of (7.7), (7.8), respectively. Thanks to the linearity of the problems (7.7), (7.8) we obtain, as for (7.17),

$$\|\varphi_\gamma^1 - \varphi_\gamma^2\|_{H_\gamma^{3/2}(\omega)} \leq \frac{C}{\gamma} \|\bar{\varphi}_\gamma^1 - \bar{\varphi}_\gamma^2\|_{H_\gamma^{3/2}(\omega)}.$$

Then there exists $\gamma_0 > 0$ such that for all $\gamma \geq \gamma_0$ the map $\bar{\varphi} \rightarrow \varphi$ has a unique fixed point, by the contraction mapping principle. The fixed point $\bar{\varphi} = \varphi$, together with the corresponding solution of (7.7), provides the solution of (7.4), (7.8), that is a solution of (6.1). As for the boundary conditions, we have already observed that (5.12), (6.1c)-(6.1f) is equivalent to (5.12), (6.1c), (7.1). The proof is complete. \square

8. Proof of Theorem 4.1

For all ε sufficiently small, problem (5.8) admits a unique solution with the regularity described in Theorem 7.2. Due to the uniform a priori estimate (6.2) we can extract a subsequence weakly convergent to functions $(\mathcal{U}, W, \varphi)$ with $(\mathcal{U}_\gamma, W_\gamma) \in H_{tan,\gamma}^1(Q_T^+) \times H_\gamma^1(Q_T^-)$ and $(q_\gamma, u_{1\gamma}, h_{1\gamma})|_{\omega_T} \in H_\gamma^{1/2}(\omega_T)$, $W_\gamma|_{\omega_T} \in H_\gamma^1(\omega_T)$ and $\varphi_\gamma \in H_\gamma^1(\omega_T)$ (we use obvious notations). Let us decompose $W = (\mathfrak{H}, \mathfrak{E})$ and perform an inverse change of unknown with respect to that of Section 5.1 to define (\mathcal{H}, E) from $(\mathfrak{H}, \mathfrak{E})$. Passing to the limit in (5.1b), (5.8)–(5.12) as $\varepsilon \rightarrow 0$ immediately gives that $(U, \mathcal{H}, \varphi)$ is a solution to (2.29), (2.23), (2.24) and $E = \mathfrak{E} = 0$. Passing to the limit in (6.2) gives the a priori estimate (4.2). The proof of Theorem 4.1 is complete.

9. Equivalence of systems (5.1b) and (6.12)

We prove the equivalence of systems (5.1b) and (6.12) for every $\vec{v} \neq 0$. This is the same as the equivalence of (5.8b) and (6.9), or (6.1b) and (6.9).

LEMMA 9.1 Assume that systems (5.1b) and (6.12) have common initial data satisfying the constraints

$$\operatorname{div} \mathfrak{h}^\varepsilon = 0, \quad \operatorname{div} \mathfrak{e}^\varepsilon = 0 \quad \text{in } \Omega^- \quad \text{for } t = 0.$$

Assuming that the corresponding Cauchy problems for (5.1b) and (6.12) have a unique classical solution on a time interval $[0, T]$, then these solutions coincide on $[0, T]$ for all ε sufficiently small.

Proof. Let us set

$$A = \hat{\eta}^{-1}(\partial_t \mathfrak{h}^\varepsilon + \varepsilon^{-1} \nabla \times \mathfrak{E}^\varepsilon), \quad B = \hat{\eta}^{-1}(\partial_t \mathfrak{e}^\varepsilon - \varepsilon^{-1} \nabla \times \mathfrak{H}^\varepsilon).$$

Then (6.12) can be written as

$$A - \varepsilon \vec{v} \times B + \frac{\vec{v}}{\partial_1 \hat{\Phi}_1} \operatorname{div} \mathfrak{h}^\varepsilon = 0, \quad B + \varepsilon \vec{v} \times A + \frac{\vec{v}}{\partial_1 \hat{\Phi}_1} \operatorname{div} \mathfrak{e}^\varepsilon = 0. \quad (9.1)$$

Taking the vector product of \vec{v} with the systems in (9.1) gives

$$\vec{v} \times A - \varepsilon \vec{v} \times (\vec{v} \times B) = 0, \quad \vec{v} \times B + \varepsilon \vec{v} \times (\vec{v} \times A) = 0, \quad (9.2)$$

that is

$$\vec{v} \times A - \varepsilon (\vec{v} \cdot B) \vec{v} + \varepsilon |\vec{v}|^2 B = 0, \quad \vec{v} \times B + \varepsilon (\vec{v} \cdot A) \vec{v} - \varepsilon |\vec{v}|^2 A = 0. \quad (9.3)$$

We take the vector product of $\varepsilon \vec{v}$ with the first system in (9.3) and get

$$\varepsilon (\vec{v} \cdot A) \vec{v} - \varepsilon |\vec{v}|^2 A + \varepsilon^2 |\vec{v}|^2 \vec{v} \times B = 0. \quad (9.4)$$

For any choice of $\vec{v} \neq 0$ we may assume that $\varepsilon |\vec{v}| \neq 1$ (this is true for ε definitely small). Then by comparison of (9.4) and the second equation in (9.3) we infer $\vec{v} \times B = 0$, and from (9.2) also $\vec{v} \times A = 0$.

Thus (6.12) may be rewritten as

$$\partial_t \mathfrak{h}^\varepsilon + \frac{1}{\varepsilon} \nabla \times \mathfrak{E}^\varepsilon + \frac{\hat{\eta} \vec{v}}{\partial_1 \hat{\Phi}_1} \operatorname{div} \mathfrak{h}^\varepsilon = 0, \quad \partial_t \mathfrak{e}^\varepsilon - \frac{1}{\varepsilon} \nabla \times \mathfrak{H}^\varepsilon + \frac{\hat{\eta} \vec{v}}{\partial_1 \hat{\Phi}_1} \operatorname{div} \mathfrak{e}^\varepsilon = 0.$$

Applying the div operator to the equations gives the transport equation

$$\partial_t u + \operatorname{div}(u \vec{a}) = 0 \quad \text{in } Q_T^-,$$

for both $u = \operatorname{div} \mathfrak{h}^\varepsilon$ and $u = \operatorname{div} \mathfrak{e}^\varepsilon$, where $\vec{a} = \hat{\eta} \vec{v} / \partial_1 \hat{\Phi}_1$. Noticing that the first component of \vec{a} vanishes at $x_1 = 0$, the transport equation doesn't need any boundary condition. As $u|_{t=0} = 0$, by a standard argument we deduce $u = 0$ for $t > 0$. This fact shows the equivalence of (5.1b) and (6.12). \square

10. Proof of Lemma 1.1

Given an even function $\chi \in C_0^\infty(\mathbb{R})$, with $\chi = 1$ on $[-1, 1]$, we define

$$\Psi(x_1, x') := \chi(x_1 \langle D \rangle) \varphi(x'), \quad (10.1)$$

where $\chi(x_1 \langle D \rangle)$ is the pseudo-differential operator with $\langle D \rangle = (1 + |D|^2)^{1/2}$ being the Fourier multiplier in the variables x' . From the definition it readily follows that $\Psi(0, x') = \varphi(x')$ for all $x' \in \mathbb{R}^2$. Moreover,

$$\partial_1 \Psi(x_1, x') = \chi'(x_1 \langle D \rangle) \langle D \rangle \varphi(x'), \quad (10.2)$$

which vanishes if $x_1 = 0$. We compute

$$\|\Psi(x_1, \cdot)\|_{H^m(\mathbb{R}^2)}^2 = \int_{\mathbb{R}^2} \langle \xi' \rangle^{2m} \chi^2(x_1 \langle \xi' \rangle) |\hat{\varphi}(\xi')|^2 d\xi',$$

where $\hat{\varphi}(\xi')$ denotes the Fourier transform in x' of φ . It follows that

$$\begin{aligned} \|\Psi\|_{L_{x_1}^2(H^m(\mathbb{R}^2))}^2 &= \int_{\mathbb{R}} \int_{\mathbb{R}^2} \langle \xi' \rangle^{2m} \chi^2(x_1 \langle \xi' \rangle) |\hat{\varphi}(\xi')|^2 d\xi' dx_1 \\ &= \int_{\mathbb{R}} \int_{\mathbb{R}^2} \langle \xi' \rangle^{2m-1} \chi^2(s) |\hat{\varphi}(\xi')|^2 d\xi' ds \leq C \|\varphi\|_{H^{m-0.5}(\mathbb{R}^2)}^2. \end{aligned}$$

In a similar way, from (10.2), we obtain

$$\begin{aligned}\|\partial_1 \Psi\|_{L^2_{x_1}(H^{m-1}(\mathbb{R}^2))}^2 &= \int_{\mathbb{R}} \int_{\mathbb{R}^2} \langle \xi' \rangle^{2m-2} |\chi'(x_1 \langle \xi' \rangle) \langle \xi' \rangle|^2 |\hat{\varphi}(\xi')|^2 d\xi' dx_1 \\ &= \int_{\mathbb{R}} \int_{\mathbb{R}^2} \langle \xi' \rangle^{2m-1} |\chi'(s)|^2 |\hat{\varphi}(\xi')|^2 d\xi' ds \leq C \|\varphi\|_{H^{m-0.5}(\mathbb{R}^2)}^2.\end{aligned}$$

Iterating the same argument yields

$$\|\partial_1^j \Psi\|_{L^2_{x_1}(H^{m-j}(\mathbb{R}^2))}^2 \leq C \|\varphi\|_{H^{m-0.5}(\mathbb{R}^2)}^2, \quad j = 0, \dots, m.$$

640 Adding over $j = 0, \dots, m$ finally gives $\Psi \in H^m(\mathbb{R}^3)$ and the continuity of the map $\varphi \mapsto \Psi$.

We now show that the cut-off function χ , and accordingly the map $\varphi \mapsto \Psi$, can be chosen to give (1.14). From (10.2) we have

$$\partial_1 \Psi(x_1, x') = (2\pi)^{-2} \int_{\mathbb{R}^2} e^{i\xi' \cdot x'} \chi'(x_1 \langle \xi' \rangle) \langle \xi' \rangle \hat{\varphi}(\xi') d\xi'.$$

By the Cauchy–Schwarz inequality and a change of variables we get

$$\begin{aligned}|\partial_1 \Psi(x)| &\leq C \|\varphi\|_{H^2(\mathbb{R}^2)} \left(\int_{\mathbb{R}^2} |\chi'(x_1 \langle \xi' \rangle)|^2 \langle \xi' \rangle^{-2} d\xi' \right)^{1/2} \\ &= C \|\varphi\|_{H^2(\mathbb{R}^2)} \left(\int_0^\infty |\chi'(x_1 \langle \rho \rangle)|^2 \langle \rho \rangle^{-2} \rho d\rho \right)^{1/2}.\end{aligned}$$

We change variables again in the integral above by setting $s = x_1 \langle \rho \rangle$. It follows that

$$|\partial_1 \Psi(x)| \leq C \|\varphi\|_{H^2(\mathbb{R}^2)} \left(\int_{x_1}^\infty |\chi'(s)|^2 \frac{x_1}{s} \frac{ds}{x_1} \right)^{1/2} \leq C \|\varphi\|_{H^2(\mathbb{R}^2)} \left(\int_1^\infty |\chi'(s)|^2 \frac{ds}{s} \right)^{1/2}. \quad (10.3)$$

Given any $M > 1$, we choose χ such that $\chi(s) = 0$ for $|s| \geq M$, and $|\chi'(s)| \leq 2/M$ for every s . Then from (10.3) one gets

$$|\partial_1 \Psi(x)| \leq \frac{C}{\sqrt{M}} \|\varphi\|_{H^2(\mathbb{R}^2)}.$$

641 Given any $\epsilon > 0$, if M is such that $C/\sqrt{M} < \epsilon$, then (1.14) immediately follows.

642 The proof of Lemma 1.2 follows from Lemma 1.1, with t as a parameter. Notice also that the
643 map $\varphi \rightarrow \Psi$, defined by (10.1), is linear and that the time regularity is conserved because, with
644 obvious notation, $\Psi(\partial_t^j \varphi) = \partial_t^j \Psi(\varphi)$. The conclusions of Lemma 1.2 follow directly.

645 11. Proof of Lemma 6.2

646 We write $\mathcal{U}_{n\gamma}^\epsilon$ on ω as

$$647 |\mathcal{U}_{n\gamma}^\epsilon|^2|_{x_1=0} = -2 \int_0^\infty \mathcal{U}_{n\gamma}^\epsilon \cdot \partial_1 \mathcal{U}_{n\gamma}^\epsilon dx_1,$$

which gives

$$\|\mathcal{U}_{n\gamma}^\varepsilon|_\omega\|_{L^2(\omega)}^2 \leq 2\|\mathcal{U}_\gamma^\varepsilon\|_{L^2(Q^+)}\|\partial_1 \mathcal{U}_{n\gamma}^\varepsilon\|_{L^2(Q^+)}. \quad (11.1)$$

Estimating the right-hand side of (11.1) with (6.17) and using the γ -homogeneity of the $H_{tan,\gamma}^1$ norm gives

$$\gamma\|\mathcal{U}_{n\gamma}^\varepsilon|_\omega\|_{L^2(\omega)}^2 \leq C \left(\|\mathcal{F}_\gamma\|_{L^2(Q^+)}^2 + \|\mathcal{U}_\gamma^\varepsilon\|_{H_{tan,\gamma}^1(Q^+)}^2 \right).$$

Thus the first part of (6.36) is proved. To show the second part of (6.36) we compute for $\ell = 0, 2, 3$,

$$\begin{aligned} \int_\omega |Z_\ell \mathcal{U}_{n\gamma}^\varepsilon|^2|_{x_1=0} dx' dt &= -2 \int_0^\infty \int_\omega Z_\ell \mathcal{U}_{n\gamma}^\varepsilon \cdot \partial_1 Z_\ell \mathcal{U}_{n\gamma}^\varepsilon dx dt \\ &= 2 \int_0^\infty \int_\omega Z_\ell^2 \mathcal{U}_{n\gamma}^\varepsilon \cdot \partial_1 \mathcal{U}_{n\gamma}^\varepsilon dx dt, \end{aligned}$$

which gives

$$\|\mathcal{U}_{n\gamma}^\varepsilon|_\omega\|_{H_\gamma^1(\omega)}^2 \leq 2\|\mathcal{U}_\gamma^\varepsilon\|_{H_{tan,\gamma}^2(Q^+)}\|\partial_1 \mathcal{U}_{n\gamma}^\varepsilon\|_{L^2(Q^+)}. \quad (11.2)$$

Interpolating between (11.1) and (11.2) gives

$$\|\mathcal{U}_{n\gamma}^\varepsilon|_\omega\|_{H_\gamma^{1/2}(\omega)}^2 \leq 2\|\mathcal{U}_\gamma^\varepsilon\|_{H_{tan,\gamma}^1(Q^+)}\|\partial_1 \mathcal{U}_{n\gamma}^\varepsilon\|_{L^2(Q^+)}.$$

Applying (6.17) eventually gives the second part of (6.36). We do the same for (6.37).

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